

*Erik Jonsson School of Engineering and Computer Science*

***Covalent Nitrogen Doping in Molecular  
Beam Epitaxy-Grown and Bulk WSe<sub>2</sub>***

**UT Dallas Author(s):**

Rafik Addou  
Christopher M. Smyth  
Ruoyu Yue  
Christopher R. Cormier  
Jiyoung Kim  
Christopher L. Hinkle  
Robert M. Wallace

**Rights:**

CC BY 4.0 (Attribution)  
©2018 The Authors

**Citation:**

Khosravi, A., R. Addou, C. M. Smyth, R. Yue, et al. 2018. "Covalent nitrogen doping in molecular beam epitaxy-grown and bulk WSe<sub>2</sub>." APL Materials 6(2), doi:10.1063/1.5002132

*This document is being made freely available by the Eugene McDermott Library of the University of Texas at Dallas with permission of the copyright owner. All rights are reserved under United States copyright law unless specified otherwise.*

## Covalent nitrogen doping in molecular beam epitaxy-grown and bulk WSe<sub>2</sub>

Ava Khosravi, Rafik Addou, Christopher M. Smyth, Ruoyu Yue,  
Christopher R. Cormier, Jiyoung Kim, Christopher L. Hinkle,  
and Robert M. Wallace<sup>a</sup>

*Department of Materials Science and Engineering, The University of Texas at Dallas,  
800 West Campbell Road, Richardson, Texas 75080, USA*

(Received 29 August 2017; accepted 8 November 2017; published online 8 January 2018)

Covalent *p*-type doping of WSe<sub>2</sub> thin films grown by molecular beam epitaxy and WSe<sub>2</sub> exfoliated from bulk crystals is achieved via remote nitrogen plasma exposure. X-ray photoelectron and Raman spectroscopies indicate covalently bonded nitrogen in the WSe<sub>2</sub> lattice as well as tunable nitrogen concentration with N<sub>2</sub> plasma exposure time. Furthermore, nitrogen incorporation induces compressive strain on the WSe<sub>2</sub> lattice after N<sub>2</sub> plasma exposure. Finally, atomic force microscopy and scanning tunneling microscopy reveal that N<sub>2</sub> plasma treatment needs to be carefully tuned to avoid any unwanted strain or surface damage. © 2018 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). <https://doi.org/10.1063/1.5002132>

Transition metal dichalcogenides (TMDs), a family of two-dimensional (2D) materials with layered structure, have gained much attention recently due to their unique material properties like atomic-scale thickness, sizeable bandgap, and potential applications in various electronic and optoelectronic devices such as tunneling field-effect transistors (FETs), p-n diodes, and photovoltaics.<sup>1–4</sup> Tungsten diselenide (WSe<sub>2</sub>) is among the most promising candidates for implementation of a TMD-based channel in next generation FETs due to its indirect-to-direct bandgap transition from bulk (~1.2 eV) to monolayer (~1.8 eV) and high *p*-type mobility up to 500 cm<sup>2</sup>/V-s at room temperature.<sup>5–7</sup> Controllable and reliable doping of TMDs is essential for realizing commercially viable TMD-based devices. Various techniques have been employed in efforts to dope WSe<sub>2</sub> such as a “mild” He or H<sub>2</sub> plasma treatment to create anion Se vacancies,<sup>8</sup> Octadecyltrichlorosilane (OTS) and hexane solution treatment,<sup>9</sup> SiN<sub>x</sub> deposition on WSe<sub>2</sub>,<sup>10</sup> potassium doping of WSe<sub>2</sub> using a K dispenser,<sup>11</sup> and Au decoration of WSe<sub>2</sub> using thermal evaporation,<sup>12</sup> none of which involve substitutional doping. Other methods such as adsorption of NO<sub>2</sub> molecules on WSe<sub>2</sub> lead to substitutional doping;<sup>13</sup> however, the formation of various WSe<sub>2-x-y</sub>O<sub>x</sub>N<sub>y</sub> species with different oxidation states was detected. A more manufacturable approach, using a remote N<sub>2</sub> plasma, has also been employed previously to introduce nitrogen into the graphene lattice (among other materials).<sup>14,15</sup> Recently, this remote N<sub>2</sub> plasma method was reported as a successful method for controllable doping of exfoliated MoS<sub>2</sub>.<sup>16</sup> In this work, we extend and compare the effect of similar N<sub>2</sub> plasma treatment on bulk WSe<sub>2</sub> and also on a few monolayers of WSe<sub>2</sub> grown by molecular beam epitaxy (MBE).

WSe<sub>2</sub> samples (thickness ≤ 7 layers) were grown via MBE on highly ordered pyrolytic graphite (HOPG) substrates provided by Momentive Performance Materials. The HOPG substrates were freshly exfoliated prior to being loaded into vacuum, followed by annealing at the 550 °C growth temperature for 1 h. Once the HOPG surface preparation was finished, the 2 h MBE growth of WSe<sub>2</sub> commenced with a W flux pressure of 7 × 10<sup>-9</sup> mbar and a Se flux pressure of 1 × 10<sup>-6</sup> mbar. Samples were slowly cooled down under a Se environment, and RHEED was performed *in situ* resulting in streaky patterns indicative of high-quality, smooth films.<sup>17</sup> The MBE samples were then removed from the growth chamber for the *ex situ* nitrogen doping experiments.

<sup>a</sup>Author to whom correspondence should be addressed: [rmwallace@utdallas.edu](mailto:rmwallace@utdallas.edu)

Bulk WSe<sub>2</sub> crystals, grown by chemical vapor transport (CVT) and purchased from HQ Graphene, were exfoliated in air using Scotch® Magic™ tape. Immediately following bulk WSe<sub>2</sub> exfoliation, both non-exfoliated MBE and exfoliated bulk WSe<sub>2</sub> samples were mounted on the same Mo-plate and loaded into an ultra-high vacuum (UHV) system described elsewhere<sup>18</sup> in less than 10 min. Samples were heated to 300 °C under UHV and held at this temperature for 2 h (base pressure  $\sim 10^{-9}$  mbar) to remove physisorbed species induced by air exposure following exfoliation. Subsequently, the samples were exposed to a N<sub>2</sub> plasma, remotely generated via 13.56 MHz Radio Frequency (RF) at a power of 100 W under 45 SCCM nitrogen gas flow, for different sets of exposure time with an operation pressure of  $\sim 7 \times 10^{-3}$  mbar. To study the surface chemistry after each treatment, the samples were transferred through a UHV transfer tube to an analysis chamber with a base pressure of  $\sim 10^{-10}$  mbar equipped with X-ray photoelectron spectroscopy (XPS) described elsewhere.<sup>18</sup> The software *AAnalyzer* was used for XPS peak analysis. Active Shirley background and Voigt line shape was employed for peak fitting.<sup>19</sup> Raman and photoluminescence (PL) spectroscopy were employed to investigate the characteristic vibrational modes and optical bandgap of monolayer and bilayer WSe<sub>2</sub> flakes and films before and after 10 min N<sub>2</sub> plasma treatment. For Raman and PL measurements, few layers WSe<sub>2</sub> samples were prepared by mechanically exfoliating the bulk WSe<sub>2</sub> crystal and transferring them to thermally grown 270 nm SiO<sub>2</sub>/Si substrates. Raman and PL spectra were acquired via 532 nm wavelength laser with a spot size of  $\sim 500$  nm and 0.22 mW power under ambient conditions. An atomic force microscope (AFM), Veeco, Model 3100 Dimension V, operating in the tapping mode under ambient conditions was also used to characterize the surface morphology *ex situ* before and after the N<sub>2</sub> plasma treatment. In a separate UHV system equipped with an *Omicron* variable temperature scanning tunneling microscope (VT-STM),<sup>20</sup> room temperature STM and scanning tunneling spectroscopy (STS) were employed to compare the surface morphology and electronic structure of the MBE WSe<sub>2</sub> films before and after the N<sub>2</sub> plasma treatment. AFM and STM images were processed with the *WSxM* software.<sup>21</sup>

Figures 1(a)–1(c) show the W 4*f*, Se 3*d*, W 5*p*<sub>3/2</sub>, and N 1*s* core level spectra obtained from as-grown MBE WSe<sub>2</sub> and as-exfoliated WSe<sub>2</sub>. Only spectral features associated with WSe<sub>2</sub> in the W 4*f* and Se 3*d* core level spectra are detected, which indicates initial WSe<sub>2</sub> surface oxidation is below detectable limits despite 10 min air exposure prior to loading into the UHV cluster tool. A small amount of adventitious carbon and oxygen are detected on the surface due to the short air exposure prior to loading into UHV as evidenced by the low-intensity chemical states detected

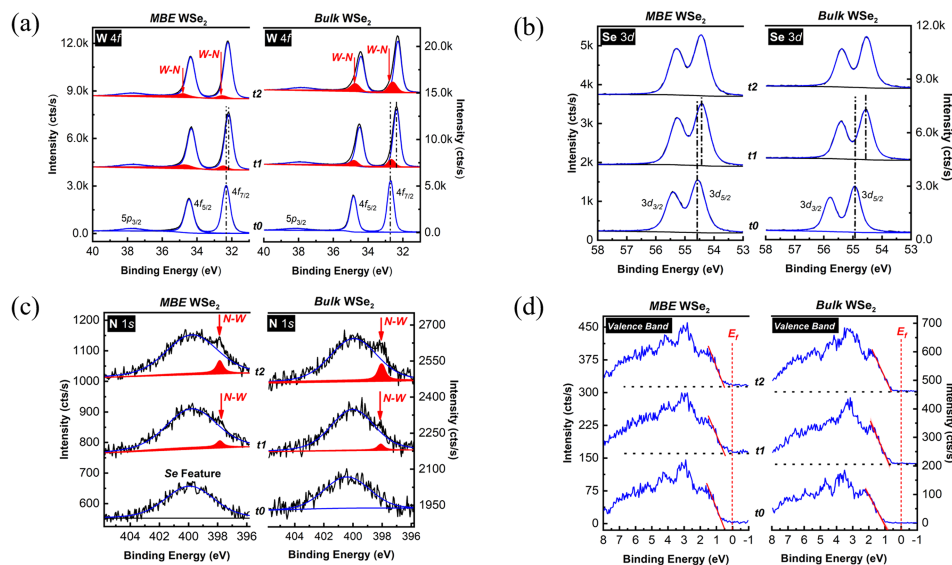


FIG. 1. [(a)–(d)] W 4*f*, Se 3*d*, W 5*p*<sub>3/2</sub>, and N 1*s* core level spectra and valence band spectra obtained from MBE and bulk, exfoliated WSe<sub>2</sub>. t0 spectra are associated with as-grown MBE WSe<sub>2</sub> and as-exfoliated WSe<sub>2</sub>, t1 with WSe<sub>2</sub> after 10 min N<sub>2</sub> exposure, and t2 with WSe<sub>2</sub> after 15 min N<sub>2</sub> exposure.

in C 1s and O 1s in Fig. S1 of the [supplementary material](#). In order to remove these undesired species and also prevent the formation of CN<sub>x</sub> during N<sub>2</sub> plasma treatment, samples were annealed at 300 °C for 2 h under UHV conditions.<sup>16</sup> It is important to note that prior to the N<sub>2</sub> plasma treatment, the N 1s core level region is convoluted with a broad Se feature at 400.0 eV, as shown in Fig. 1(c). Nitrogen remains below the XPS detection limit following N<sub>2</sub> plasma treatment for less than 10 min. For exposure times ≥10 min, the binding energy of the chemical state detected in the N 1s core level (397.9 eV) is consistent with previous reports of a transition metal-nitrogen bond.<sup>22,23</sup> An additional chemical state is detected in the W 4f core level obtained from both MBE and bulk WSe<sub>2</sub> at ~0.3 eV higher binding energy to that of W–Se after 10 min N<sub>2</sub> plasma exposure. The binding energy of the chemical state detected in the N 1s core level, in conjunction with the additional chemical state detected in the W 4f core level after 10 min N<sub>2</sub> plasma exposure, suggests the formation of covalent W–N bonds during the N<sub>2</sub> plasma treatment. Previous reports on W–N chemical states have shown similar binding energies to those reported here.<sup>22–24</sup> The detection of chemical states consistent with covalent W–N bonding in the W 4f and N 1s core level spectra obtained from both MBE and bulk WSe<sub>2</sub> samples suggests nitrogen incorporation into WSe<sub>2</sub> through Se substitution.

No additional chemical state was detected in the Se 3d core level aside from that of WSe<sub>2</sub> through all N<sub>2</sub> plasma exposure times investigated in this work, which indicates Se–N bond formation, if any, is below the limit of detection by XPS. After 20, 30, and 60 min exposure time, the intensity of the chemical state detected in the N 1s core level obtained from both MBE and bulk WSe<sub>2</sub> gradually increases (Fig. S2 of the [supplementary material](#)). Increasing the N<sub>2</sub> plasma exposure time results in an increased nitrogen to tungsten ratio and corresponding decrease in the selenium to tungsten ratio in WN<sub>x</sub>Se<sub>2–x</sub>. After 60 min N<sub>2</sub> plasma exposure,  $x = 0.17$  for MBE WSe<sub>2</sub> and  $x = 0.16$  for bulk WSe<sub>2</sub> as calculated from the XPS data. It should be noted that the W–N bond represents the only nitrogen species detected in either sample after N<sub>2</sub> plasma exposure of any duration. According to angle resolved XPS (ARXPS) performed at a bulk sensitive angle (75°) *versus* a surface sensitive angle (45°), the remote N<sub>2</sub> plasma exposure introduces nitrogen to the surface of WSe<sub>2</sub> samples (Fig. S3 of the [supplementary material](#)). Taken together, the ARXPS and detected surface chemical states are consistent with Se replacement by N from the remote plasma process, analogous to that reported for MoS<sub>2</sub>.

The XPS spectra also shows a clear shift to lower binding energy by the same amount for all core levels, indicative of a Fermi level shift. The valence band maximum (VBM) was extrapolated from linear fits of the valence band spectra obtained throughout the experiment [Fig. 1(d)]. A valence band edge shift to lower binding energy (toward the Fermi level) is an indicator of *p*-type doping induced by N<sub>2</sub> plasma treatment.<sup>9</sup> *P*-type doping behavior upon Se substitution by nitrogen is expected considering the electronegativity of nitrogen comparing to selenium.<sup>25</sup> After 10 min exposure to N<sub>2</sub> plasma, the valence band edge shifts detected from MBE WSe<sub>2</sub> (0.2 eV) and bulk WSe<sub>2</sub> (0.5 eV) were consistent with the shifts in the W 4f and Se 3d core levels. Fermi level shifts toward the valence band edge (*p*-type) were larger in bulk WSe<sub>2</sub> compared to MBE WSe<sub>2</sub> in all sets of samples investigated in this work. The difference in the Fermi level shift detected between MBE and bulk WSe<sub>2</sub> samples could be explained by the initial VBM position of the respective samples. The as-grown MBE film is initially *p*-type as indicated by the VBM located 0.7 eV below the Fermi level while the VBM initially detected from bulk WSe<sub>2</sub> is 1.2 eV below E<sub>F</sub>. These differences could be due to extrinsic impurities in the bulk crystals.<sup>6</sup>

In comparison with the study on MoS<sub>2</sub>,<sup>16</sup> nitrogen incorporation into WSe<sub>2</sub> occurs at a significantly slower rate as nitrogen is detectable after 10 min N<sub>2</sub> plasma exposure compared to just 2 min in the case of the same treatment applied to MoS<sub>2</sub>. In addition, after 60 min N<sub>2</sub> plasma exposure, the nitrogen to tungsten ratio in WN<sub>x</sub>Se<sub>2–x</sub> is much lower than the nitrogen to molybdenum ratio in MoN<sub>x</sub>S<sub>2–x</sub> obtained after the same duration of plasma exposure.

Similar to the covalent doping of MoS<sub>2</sub> by N<sub>2</sub> plasma,<sup>16</sup> the incorporation of nitrogen into WSe<sub>2</sub> samples and formation of the W–N bond could result in compressive strain due to a smaller atomic radius of nitrogen compared to that of selenium. Raman spectroscopy is known to be a powerful tool to reveal the effect of doping on the sample structure. Raman spectroscopy was employed to qualitatively investigate the effects of N<sub>2</sub> plasma on lattice strain in exfoliated monolayer (ML) and bilayer (BL)



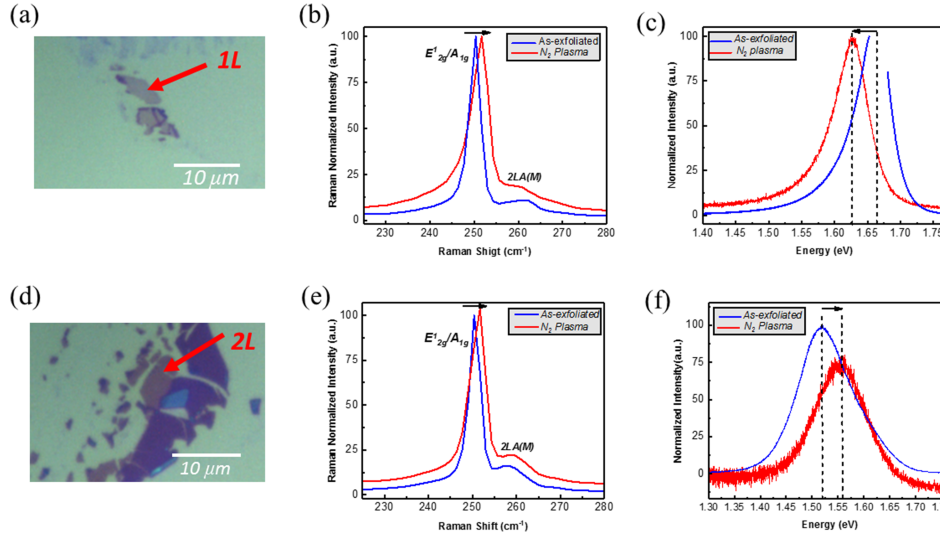


FIG. 2. Optical images of exfoliated (a) monolayer and (d) bilayer WSe<sub>2</sub> flakes. Raman spectra obtained from (b) monolayer and (e) bilayer WSe<sub>2</sub>. PL spectra obtained from the same (c) monolayer and (f) bilayer WSe<sub>2</sub> flakes.

WSe<sub>2</sub> flakes. Figures 2(a) and 2(d) show optical images of ML and BL WSe<sub>2</sub> flakes, respectively. The separation between the peak convoluted with degenerate  $E'_{2g}/A_{1g}$  vibrational modes and the higher wavenumber peak corresponding with the second order Raman mode due to LA phonons at the M point in the Brillouin zone [2LA(M)]<sup>26</sup> is a fingerprint of number of layers of WSe<sub>2</sub>.<sup>26,27</sup> The separation between the degenerate  $E'_{2g}/A_{1g}$  peak and 2LA(M) peak detected from the flakes shown in Figs. 2(a) and 2(d) is  $\sim 11 \text{ cm}^{-1}$  and  $\sim 6 \text{ cm}^{-1}$ , consistent with ML and BL WSe<sub>2</sub>, respectively [Figs. 2(b) and 2(e)]. After 10 min N<sub>2</sub> plasma exposure, the  $E'_{2g}/A_{1g}$  peak blue shifts  $2.3 \text{ cm}^{-1}$  and  $1.4 \text{ cm}^{-1}$  for the monolayer and bilayer samples, respectively. The blue shift in the  $E'_{2g}/A_{1g}$  peak indicates compressive strain from a reduction in the lattice constant<sup>28</sup> and also *p*-type doping. *P*-type doping has also been correlated with an intensity increase of the  $A_{1g}$  phonon.<sup>9</sup>

To examine the optical bandgap, PL spectroscopy was employed on the same exfoliated WSe<sub>2</sub> flakes as were probed with Raman spectroscopy. The peak position at 1.66 eV detected by PL from the ML WSe<sub>2</sub> flake [Fig. 2(c)] corresponds with the direct bandgap of monolayer WSe<sub>2</sub>.<sup>28</sup> Monolayer WSe<sub>2</sub>, of course, exhibits the highest PL intensity among ML, BL, Tri-layer, and quad-layer WSe<sub>2</sub> flakes,<sup>26</sup> with a maximum intensity over the limit of the detector employed in this work. Figure 2(f) shows the PL spectrum obtained from bilayer WSe<sub>2</sub> with peak intensity at 1.52 eV, consistent with a previous report.<sup>28</sup> A red shift in PL peak is detected after 10 min N<sub>2</sub> plasma exposure in the case of monolayer WSe<sub>2</sub> [Fig. 2(c)], consistent with a change in band structure due to induced strain caused by the covalently incorporated nitrogen.<sup>28</sup> Desai *et al.*<sup>28</sup> report a crossover from indirect to direct bandgap in bilayer WSe<sub>2</sub> is practically possible by applying strain due to the small difference between the direct and indirect bandgaps of bilayer WSe<sub>2</sub>. As strain increases, a second peak, corresponding with the direct bandgap, appears at higher energy in the PL spectrum and increases in intensity until the indirect bandgap peak is no longer resolvable. Further increase in strain results in a red shift of the direct bandgap peak of bilayer WSe<sub>2</sub>. In this work, the PL spectra detected from bilayer WSe<sub>2</sub> after 10 min N<sub>2</sub> plasma exposure exhibits characteristics reflecting a lesser magnitude of strain than that corresponding with a PL spectrum dominated by the direct bandgap peak as observed by Desai *et al.* Both Raman and PL provide evidence of successful nitrogen incorporation at the surface of mono and bilayer WSe<sub>2</sub> flakes resulting in *p*-type doping and inducing strain.

This work demonstrates covalent *p*-type doping of both MBE and bulk WSe<sub>2</sub> via a remote N<sub>2</sub> plasma treatment. However, surface roughening and/or etching could be an undesirable side effect of the N<sub>2</sub> plasma exposure to WSe<sub>2</sub> in a device fabrication procedure.<sup>23</sup> The WSe<sub>2</sub> surface topography after N<sub>2</sub> plasma exposure was evaluated by AFM and STM. Figure 3(a) shows the flat

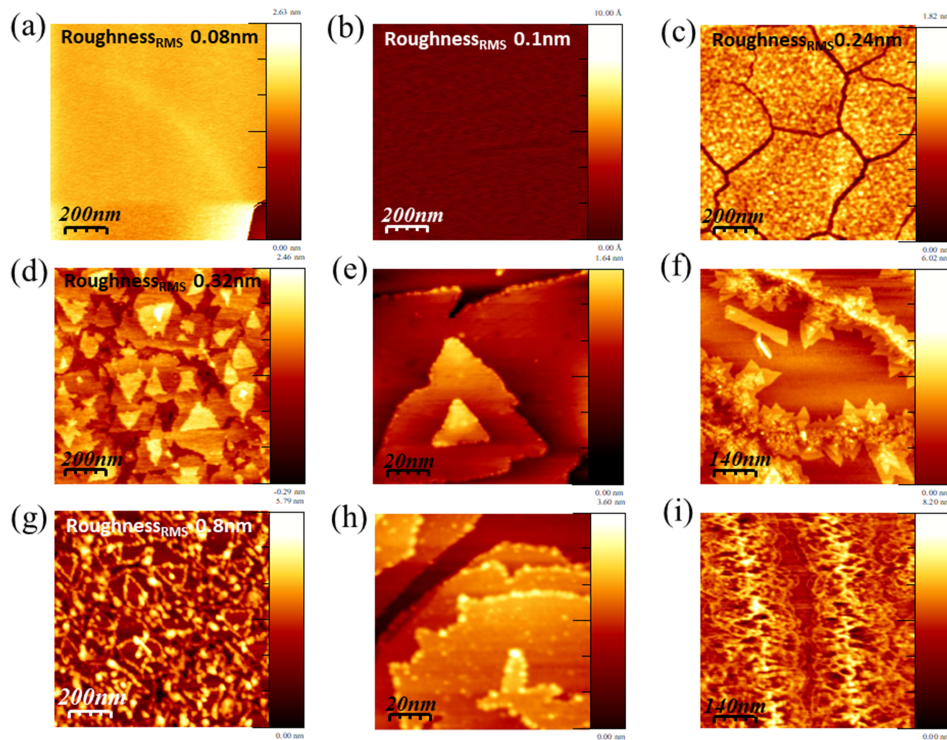


FIG. 3. AFM images obtained from exfoliated WSe<sub>2</sub> (a) before treatment, (b) after 15 min, and (c) after 60 min N<sub>2</sub> plasma exposure. AFM images of MBE WSe<sub>2</sub> (d) before treatment and (g) after 15 min N<sub>2</sub> plasma exposure. STM images obtained from the same MBE WSe<sub>2</sub> [(e) and (f)] before N<sub>2</sub> plasma treatment with the tip bias of 1.8 V and tunneling current of 300 pA and [(h) and (i)] after 30 min treatment with tip bias of 2 V, 1.8 V, and tunneling current of 200 pA.

surface of bulk WSe<sub>2</sub> before N<sub>2</sub> plasma treatment. After N<sub>2</sub> plasma exposure for 15 min, there are no obvious topographical changes [Fig. 3(b)] in the surface of bulk WSe<sub>2</sub> consistent with an analogous N<sub>2</sub> plasma treatment on bulk MoS<sub>2</sub> for short exposure times (15 min).<sup>16</sup> After 60 min of the N<sub>2</sub> plasma treatment, significant changes in the surface structure appear in the form of cracks with depth of ~0.7 nm, the height of one layer of WSe<sub>2</sub> [Fig. 3(c)]. Well-defined, mono and few-layer triangular WSe<sub>2</sub> domains with sharp (oxidation-free) edges are observed on the as-grown MBE WSe<sub>2</sub> surface [Fig. 3(d)]. After 15 min N<sub>2</sub> plasma exposure, the AFM image in Fig. 3(g) indicates a strong edge decoration on all WSe<sub>2</sub> domains suggesting that nitrogen preferentially reacts with the edges.

The contrast at the as-grown WSe<sub>2</sub> domain edges is comparable with that of the terraces suggesting the surface potential is similar across the probed region. However, N<sub>2</sub> plasma exposure for 30 min results in a sharp contrast between the terraces (dark) and domain edges (bright, difference in the electron density, and/or the height), which corresponds to significant change in surface potential in the vicinity of the domain edges [Figs. 3(e) and 3(h)]. Edge roughening observed in AFM images concomitant with the edge localized contrast variation detected with STM suggests that nitrogen preferentially nucleates at WSe<sub>2</sub> edges, which are known to be more reactive than the WSe<sub>2</sub> basal plane.<sup>29</sup> The contrast variation observed in STM images could relate to the metallic character of the W–N bond, and the tunneling current is higher in these areas compared with the adjacent semi-conducting WSe<sub>2</sub>. Figures 3(f) and 3(i) present STM images obtained from the MBE WSe<sub>2</sub> film, showing a high density of WSe<sub>2</sub> domains at HOPG step edges with small domain size as previously reported.<sup>17</sup> Following 30 min remote N<sub>2</sub> plasma treatment, there is a noticeable degradation of the surface as evidenced by a large increase in surface roughness [Fig. S4(b) of the [supplementary material](#)]. Damage is observed across the entire surface of smaller MBE WSe<sub>2</sub> domains (Fig. S4 of the [supplementary material](#)). In contrast, large MBE WSe<sub>2</sub> domains appear to only be affected on the edges [Fig. 3(h)]. Both AFM and STM results suggest that nitridation occurs preferentially at

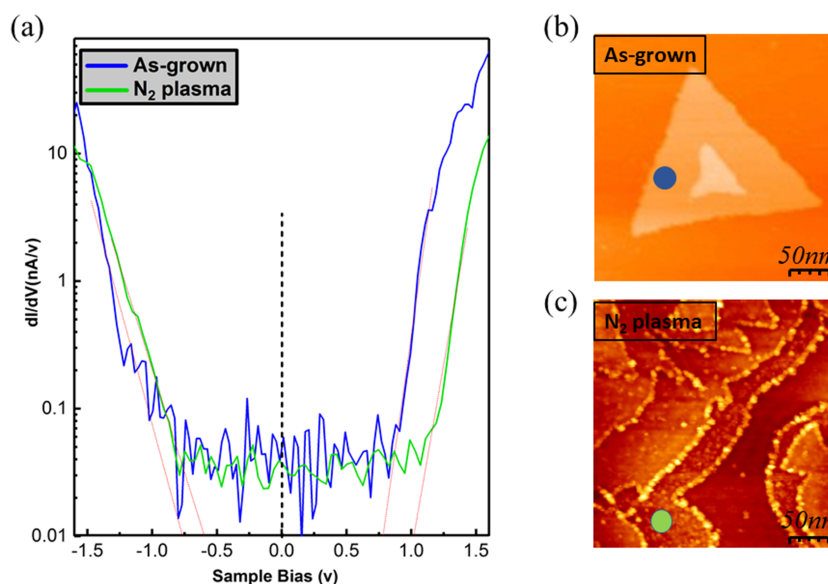


FIG. 4. (a) Comparison of scanning tunneling spectroscopy spectra measured (b) before N<sub>2</sub> plasma treatment with tip bias of 2 V and tunneling current of 300 pA (blue) and (c) after 30 min N<sub>2</sub> plasma treatment with the tip bias of 1.5 V and tunneling current of 500 pA (red). The valence band maximum (VBM) and conduction band minimum (CBM) was measured on as-grown MBE WSe<sub>2</sub> monolayer at  $-0.98$  eV and  $0.85$  eV, respectively. Following 30 min N<sub>2</sub> plasma treatment, the VBM (CBM) was measured at  $1.14$  eV ( $-0.67$  eV).

the domain edges prior to affecting the WSe<sub>2</sub> terraces. Figure 4 displays the STS spectra obtained from an MBE WSe<sub>2</sub> monolayer prior to and following N<sub>2</sub> plasma treatment. A  $1.83$  eV bandgap is measured from as-grown monolayer WSe<sub>2</sub> with no significant changes induced by the N<sub>2</sub> plasma treatment ( $1.81$  eV after 30 min exposure). The Fermi level moves toward the valence band maximum (VBM) by  $0.3$  eV in agreement with the Fermi level shift observed by XPS. This is further evidence that N<sub>2</sub> plasma exposure induces *p*-type doping in MBE WSe<sub>2</sub>.

In summary, the effects of N<sub>2</sub> plasma exposure on the electronic and structural properties of MBE thin films and bulk WSe<sub>2</sub> are evaluated using XPS, Raman, PL, AFM, and STM. The appearance of an additional chemical state in the W  $4f$  core level in conjunction with increased nitrogen concentration above the limit of XPS detection following N<sub>2</sub> plasma exposure indicates that nitrogen forms a covalent bond with tungsten. Shifts detected in characteristic WSe<sub>2</sub> vibrational modes and the PL peak after N<sub>2</sub> plasma exposure suggest that incorporating nitrogen into the WSe<sub>2</sub> lattice in the form of covalent W–N bonds induces strain in both MBE and bulk WSe<sub>2</sub>. In addition, AFM and STM images indicate that N<sub>2</sub> plasma exposure damages the WSe<sub>2</sub> surface after longer exposure times eventually resulting in the formation of cracks in the case of bulk WSe<sub>2</sub>, while the remote N<sub>2</sub> plasma process employed here is more destructive to few-layer MBE WSe<sub>2</sub> compared to bulk WSe<sub>2</sub>. Tuning the duration of N<sub>2</sub> plasma exposure offers control over the amount of strain induced in the WSe<sub>2</sub> structure, both bulk and MBE, and can pave the ground for engineering the bandgap and electronic properties of WSe<sub>2</sub>.<sup>30</sup>

See [supplementary material](#) for the O  $1s$  and C  $1s$  core levels and Valence Band Maximum of MBE and exfoliated WSe<sub>2</sub> prior to load them into UHV system and after annealing, the sequential N<sub>2</sub> plasma exposure and gradual intensity increase of W–N peak is demonstrated, the angle resolved XPS in bulk sensitive angle from both samples is provided, XPS spectra from HOPG as a substrate after N<sub>2</sub> plasma exposure, and STM image from MBE WSe<sub>2</sub> after 30 min N<sub>2</sub> exposure is included.

This work is supported in part by the Center for Low Energy Systems Technology (LEAST), one of six centers supported by the STARnet phase of the Focus Center Research Program (FCRP), a Semiconductor Research Corporation program sponsored by MARCO and DARPA. Also, by the Southwest Academy on Nanoelectronics (SWAN) sponsored by the Nano-electronic Research

Initiative and NIST and in part by NSF Award No. 1407765 under the US/Ireland UNITE collaboration. The authors acknowledge Dr. Lanxia Cheng for useful discussions on the Raman and PL study.

- <sup>1</sup> W. Liu, J. Kang, D. Sarkar, Y. Khatami, D. Jena, and K. Banerjee, *Nano Lett.* **13**(5), 1983–1990 (2013).
- <sup>2</sup> H. Fang, S. Chuang, T. C. Chang, K. Takei, T. Takahashi, and A. Javey, *Nano Lett.* **12**(7), 3788–3792 (2012).
- <sup>3</sup> B. W. Baugher, H. O. Churchill, Y. Yang, and P. Jarillo-Herrero, *Nat. Nanotechnol.* **9**(4), 262–267 (2014).
- <sup>4</sup> A. Pospischil, M. M. Furchi, and T. Mueller, *Nat. Nanotechnol.* **9**(4), 257–261 (2014).
- <sup>5</sup> M. Deshpande, G. Solanki, and M. Agarwal, *Mater. Lett.* **43**(1), 66–72 (2000).
- <sup>6</sup> R. Addou and R. M. Wallace, *ACS Appl. Mater. Interfaces* **8**(39), 26400–26406 (2016).
- <sup>7</sup> V. Podzorov, M. Gershenson, C. Kloc, R. Zeis, and E. Bucher, *Appl. Phys. Lett.* **84**(17), 3301–3303 (2004).
- <sup>8</sup> M. Tosun, L. Chan, M. Amani, T. Roy, G. H. Ahn, P. Taheri, C. Carraro, J. W. Ager, R. Maboudian, and A. Javey, *ACS Nano* **10**(7), 6853–6860 (2016).
- <sup>9</sup> D.-H. Kang, J. Shim, S. K. Jang, J. Jeon, M. H. Jeon, G. Y. Yeom, W.-S. Jung, Y. H. Jang, S. Lee, and J.-H. Park, *ACS Nano* **9**(2), 1099–1107 (2015).
- <sup>10</sup> K. Chen, D. Kiriya, M. Hettick, M. Tosun, T.-J. Ha, S. R. Madhupathy, S. Desai, A. Sachid, and A. Javey, *APL Mater.* **2**(9), 092504 (2014).
- <sup>11</sup> H. Fang, M. Tosun, G. Seol, T. C. Chang, K. Takei, J. Guo, and A. Javey, *Nano Lett.* **13**(5), 1991–1995 (2013).
- <sup>12</sup> C.-H. Chen, C.-L. Wu, J. Pu, M.-H. Chiu, P. Kumar, T. Takenobu, and L.-J. Li, *2D Mater.* **1**(3), 034001 (2014).
- <sup>13</sup> P. Zhao, D. Kiriya, A. Azcatl, C. Zhang, M. Tosun, Y.-S. Liu, M. Hettick, J. S. Kang, S. McDonnell, and K. C. Santhosh, *ACS Nano* **8**(10), 10808–10814 (2014).
- <sup>14</sup> K. Akada, T.-O. Terasawa, G. Imamura, S. Obata, and K. Saiki, *Appl. Phys. Lett.* **104**(13), 131602 (2014).
- <sup>15</sup> A. Talbi, Z. Benamara, B. Akkal, B. Gruzza, L. Bideux, C. Robert, C. Varenne, and N. Chami, *Mater. Sci. Eng.: A* **437**(2), 254–258 (2006).
- <sup>16</sup> A. Azcatl, X. Qin, A. Prakash, C. Zhang, L. Cheng, Q. Wang, N. Lu, M. J. Kim, J. Kim, and K. Cho, *Nano Lett.* **16**(9), 5437–5443 (2016).
- <sup>17</sup> R. Yue, Y. Nie, L. Walsh, R. Addou, C. Liang, N. Lu, A. Barton, H. Zhu, Z. Che, and D. Barrera, *2D Mater.* **4**, 045019 (2017).
- <sup>18</sup> R. M. Wallace, *ECS Trans.* **16**(5), 255–271 (2008).
- <sup>19</sup> A. Herrera-Gomez, A. Hegedus, and P. Meissner, *Appl. Phys. Lett.* **81**(6), 1014–1016 (2002).
- <sup>20</sup> R. Addou, L. Colombo, and R. M. Wallace, *ACS Appl. Mater. Interfaces* **7**(22), 11921–11929 (2015).
- <sup>21</sup> I. Horcas, R. Fernández, J. Gomez-Rodriguez, J. Colchero, J. Gómez-Herrero, and A. Baro, *Rev. Sci. Instrum.* **78**(1), 013705 (2007).
- <sup>22</sup> K. Schmid, A. Manhard, C. Linsmeier, A. Wiltner, T. Schwarz-Selinger, W. Jacob, and S. Mändl, *Nucl. Fusion* **50**(2), 025006 (2010).
- <sup>23</sup> D. Alegre, T. Acsente, A. Martin-Rojo, E. Oyarzabal, F. Tabares, G. Dinescu, G. De Temmerman, R. Birjega, C. Logofatu, and J. Kovac, *Rom. Rep. Phys.* **67**(2), 532–546 (2015).
- <sup>24</sup> M. Nagai, T. Suda, K. Oshikawa, N. Hirano, and S. Omi, *Catal. Today* **50**(1), 29–37 (1999).
- <sup>25</sup> Y. Zhang, *Inorg. Chem.* **21**(11), 3886–3889 (1982).
- <sup>26</sup> W. Zhao, Z. Ghorannevis, K. K. Amara, J. R. Pang, M. Toh, X. Zhang, C. Kloc, P. H. Tan, and G. Eda, *Nanoscale* **5**(20), 9677–9683 (2013).
- <sup>27</sup> H. Sahin, S. Tongay, S. Horzum, W. Fan, J. Zhou, J. Li, J. Wu, and F. Peeters, *Phys. Rev. B* **87**(16), 165409 (2013).
- <sup>28</sup> S. B. Desai, G. Seol, J. S. Kang, H. Fang, C. Battaglia, R. Kapadia, J. W. Ager, J. Guo, and A. Javey, *Nano Lett.* **14**(8), 4592–4597 (2014).
- <sup>29</sup> R. C. Longo, R. Addou, K. Santosh, J.-Y. Noh, C. M. Smyth, D. Barrera, C. Zhang, J. W. Hsu, R. M. Wallace, and K. Cho, *2D Mater.* **4**(2), 025050 (2017).
- <sup>30</sup> P. Johari and V. B. Shenoy, *ACS Nano* **5**(7), 5903–5908 (2011).