

School of Natural Sciences and Mathematics

***Search for Heavy Higgs Bosons A / H Decaying to a Top Quark
Pair in pp Collisions at $\sqrt{s} = 8$ TeV with the ATLAS Detector***

UT Dallas Author(s):

Joseph M. Izen
Bernhard Meirose
Kendall Reeves

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Search for Heavy Higgs Bosons A/H Decaying to a Top Quark Pair in pp Collisions at $\sqrt{s} = 8$ TeV with the ATLAS Detector

M. Aaboud *et al.**

(ATLAS Collaboration)

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A search for heavy pseudoscalar (A) and scalar (H) Higgs bosons decaying into a top quark pair ($t\bar{t}$) has been performed with 20.3 fb^{-1} of proton-proton collision data collected by the ATLAS experiment at the Large Hadron Collider at a center-of-mass energy $\sqrt{s} = 8$ TeV. Interference effects between the signal process and standard model $t\bar{t}$ production, which are expected to distort the signal shape from a single peak to a peak-dip structure, are taken into account. No significant deviation from the standard model prediction is observed in the $t\bar{t}$ invariant mass spectrum in final states with an electron or muon, large missing transverse momentum, and at least four jets. The results are interpreted within the context of a type-II two-Higgs-doublet model. Exclusion limits on the signal strength are derived as a function of the mass $m_{A/H}$ and the ratio of the vacuum expectation values of the two Higgs fields, $\tan\beta$, for $m_{A/H} > 500$ GeV.

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Introduction.—The production of new particles at the Large Hadron Collider (LHC) with masses close to the TeV scale is predicted by many models of physics beyond the standard model (SM). In this Letter, a search for massive pseudoscalar and scalar resonances decaying into a top-antitop quark pair ($t\bar{t}$) is presented. It is the first search in this final state to take into account the significant interference between the signal and the background from SM $t\bar{t}$ production. The search is conducted on a sample of pp collision data with an integrated luminosity of 20.3 fb^{-1} at a center-of-mass energy $\sqrt{s} = 8$ TeV, collected with the ATLAS detector [1].

New pseudoscalar (A) and scalar (H) states coupling strongly to $t\bar{t}$ are predicted by a class of models in which the Higgs sector is extended to include a second Higgs doublet, the two-Higgs-doublet models (2HDMs) [2]. These models are motivated by many theories beyond the SM, such as supersymmetry [3–8] and axion models [9]. In 2HDMs of type II [2], such as the minimal supersymmetric standard model (MSSM) [10–14], these states decay predominantly into $t\bar{t}$ pairs if $m_{A/H} \geq 500$ GeV and the ratio of the vacuum expectation values of the two Higgs fields, $\tan\beta$, is small ($\tan\beta \lesssim 3$).

To date, this parameter region has not been probed directly by searches in other final states [15–20] or by previous searches for $t\bar{t}$ resonances [21–25]. The latter,

which aim to identify resonant excesses in the $t\bar{t}$ invariant mass ($m_{t\bar{t}}$) spectrum, have a reduced sensitivity to 2HDM signatures as they do not take into account interference effects between the signal and the dominant background from SM $t\bar{t}$ production. These are significant for (pseudo) scalar Higgs bosons with masses above the $t\bar{t}$ production threshold where the interference between the gluon-gluon (gg) initiated loop production and the irreducible background from SM $t\bar{t}$ production yields a non-negligible imaginary term in the amplitude, which at the LHC is dominated by $gg \rightarrow t\bar{t}$ production [26–31]. As a result of the interference, the signal shape is distorted from a Breit-Wigner peak to a peak-dip structure.

The results of the search are interpreted in a CP -conserving type-II 2HDM with a softly broken Z_2 symmetry [32]. The lighter of the two neutral CP -even states, h , is assumed to be the Higgs boson discovered at a mass of $m_h = 125$ GeV [33,34] with couplings as predicted by the SM. This corresponds to the condition $\sin(\alpha - \beta) = 1$, referred to as the *alignment limit*, where α denotes the mixing angle between the two CP -even states. The parameter m_{12} of the Z_2 breaking term of the potential is taken to be $m_{12}^2 = m_A^2 \tan\beta / (1 + \tan^2\beta)$. In this model, the production cross sections and widths of A and H , as well as the signal shape, are determined by $\tan\beta$ and the masses m_A and m_H . The search results are derived assuming mass degeneracy, $m_H = m_A$, such that both processes contribute to the $m_{t\bar{t}}$ spectrum, a scenario motivated, for example, by the MSSM [32]. We also consider two scenarios in which only the interference pattern of either A or H appears in the $m_{t\bar{t}}$ spectrum [35].

Data and Monte Carlo samples.—This analysis closely follows the resolved-topology analysis in Ref. [22]. Events with signatures compatible with $t\bar{t} \rightarrow W^+ b W^- \bar{b}$, with one

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W boson decaying hadronically and the other leptonically, the lepton-plus-jets channel ($\ell + \text{jets}$, $\ell = e, \mu$), were collected using single-electron and single-muon triggers. The trigger efficiency is constant in the transverse momentum (p_T) of leptons with $p_T > 25$ GeV [36,37]. The dominant background arises from SM $t\bar{t}$ production, followed by a contribution from $W + \text{jets}$ processes. Data-driven techniques were used to normalize the $W + \text{jets}$ background contribution and to estimate the background from multijet events. All other background processes were estimated using Monte Carlo (MC) simulation. The background estimates for all processes are identical to those in Ref. [22].

The signal process $gg \rightarrow A/H \rightarrow t\bar{t}$, including the decays of the top quarks and resulting W bosons, was simulated using MADGRAPH5_aMC@NLO [38] v2.3.3 with the model of Ref. [39], which implements the A/H production through loop-induced gluon-gluon fusion with loop contributions from top and bottom quarks at leading order (LO) in QCD. The CT10 set [40] of parton distribution functions (PDFs) was used and the renormalization and factorization scales were set to $\sqrt{\sum_{\text{decay products}}(p_T^2 + m^2)}$.

For the statistical interpretation, the $t\bar{t}$ invariant mass distributions in the signal regions in data were compared to a combination of the expected distributions from all background processes B , the pure signal process S , and the signal-plus-interference component $S + I$ for a given signal hypothesis, as illustrated in Eq. (1) below. The most reliable description of the $t\bar{t}$ background [41] is obtained at next-to-leading order (NLO) with POWHEG-BOX [42–45] + PYTHIA6[46]. Therefore, the $S + I$ contribution was modeled separately from this background process by modifying the MADGRAPH5_aMC@NLO software to remove the pure SM $t\bar{t}$ process to yield only the $S + I$ contribution on an event-by-event basis. The nominal $t\bar{t}$ background prediction in $m_{t\bar{t}}$ is in good agreement with that obtained from MADGRAPH5_aMC@NLO in all signal regions. The $S + I$ events obtained with the modified software can have positive or negative weights. Figure 1 shows the $t\bar{t}$ invariant mass distributions for the S and $S + I$ components in a model with $\tan\beta = 0.68$ and a pseudoscalar of mass $m_A = 500$ GeV. The $S + I$ component exhibits a peak-dip structure with the minimum around $m_{A/H}$ for all signal hypotheses studied in this search. The width of both the S and $S + I$ distribution decreases with increasing $\tan\beta$.

The $S + I$ distributions from the modified MADGRAPH5_aMC@NLO software were validated against those from the unmodified program. The latter were obtained by generating a large inclusive sample $S + I + B_{t\bar{t}}$ for a given parameter point and a LO SM $t\bar{t}$ background $B_{t\bar{t}}$ sample with the same generator settings. The difference between the resulting two $m_{t\bar{t}}$ distributions corresponds to the $S + I$ component, which agrees with

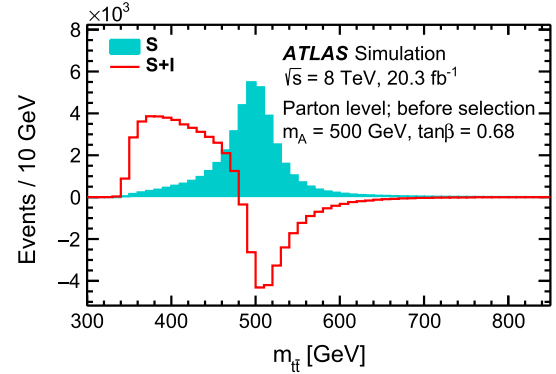


FIG. 1. Distributions of the invariant mass of the $t\bar{t}$ pair from the decay of a pseudoscalar A of mass $m_A = 500$ GeV before the emission of final-state radiation and before the parton shower for the pure resonance S (filled) and signal + interference contribution $S + I$ (unfilled). Events from all $t\bar{t}$ decay modes are included.

that obtained with the modified software within 0.4% across the whole spectrum. The difference is taken as a systematic uncertainty in $S + I$.

PYTHIA6 with the Perugia 2011c set of tuned parameters [47] was used to model the parton shower and hadronization for all S and $S + I$ samples and the stable particles obtained after hadronization were passed through the ATLAS fast detector simulation [48]. The effects of additional collisions within the same or nearby bunch crossings were simulated by overlaying additional pp collisions, simulated with PYTHIA v8.1 [49], on each event. Correction factors were applied to adjust the trigger and selection efficiencies in simulated events to those measured in data. The S and $S + I$ samples with this setup were generated separately for pseudoscalar and scalar Higgs bosons.

Event samples for both the S and $S + I$ components for different values of $(m_{A/H}, \tan\beta)$ were obtained from signal samples S after the detector simulation by applying an event-by-event reweighting. This reweighting substantially reduces the computing time required. The weight is the ratio of the MADGRAPH5_aMC@NLO matrix elements, calculated from the four-momenta of the incoming gluons and outgoing top quarks of the generated event with the new and the old values of $(m_{A/H}, \tan\beta)$, respectively. All $S + I$ and a small number of S samples were obtained through reweighting. Signal hypotheses with $m_{A/H} < 500$ GeV were not considered as they require an accurate modeling of the Higgs boson decay into virtual top quarks and the implementation of higher-order corrections that are not available in the MADGRAPH5_aMC@NLO model. The requirement $\tan\beta \geq 0.4$ was imposed to ensure the perturbativity of the top-quark Yukawa coupling [2].

Correction factors K_S were applied to normalize the generated signal (S) cross section to the value calculated at partial next-to-next-to-leading-order (NNLO) precision in

QCD [50–52]. The correction factor for the interference component I is $K_I = \sqrt{K_S \times K_B}$, as suggested in Ref. [53], where $K_B = 1.87$ is the correction factor to normalize the total cross section of the SM $t\bar{t}$ background generated at LO with MADGRAPH to the cross section calculated at NNLO accuracy in the strong coupling constant α_s , including resummation of next-to-next-to-leading-logarithmic soft gluon terms. The values of K_S range between two and three for the tested signal hypotheses.

Event selection.—The event selection criteria for the signal regions provide a high selection efficiency for $t\bar{t}$ events. Only events with a *resolved topology*, in which the three jets from the hadronically decaying top quark are well separated in the detector, are selected. This is the most efficient selection strategy for signal hypotheses with $m_{A/H} < 800$ GeV. Events with a *merged topology*, in which the top quark is reconstructed as a single jet, are not considered. The event reconstruction and selection criteria are identical to those in Ref. [22] except that events that would satisfy the criteria for both topologies are classified as “resolved” instead of “merged.”

Events are required to contain exactly one isolated electron [54] or muon [55] with $p_T > 25$ GeV and pseudorapidity $|\eta| < 2.5$ [56]. Events must have large missing transverse momentum, $E_T^{\text{miss}} > 20$ GeV, computed as the magnitude of the negative vector sum of lepton and jet transverse momenta [57]. In addition, $E_T^{\text{miss}} + m_T^W > 60$ GeV is required to further suppress the contribution from multijet events, where m_T^W is the lepton– E_T^{miss} transverse mass [22]. Events must contain at least four hadronic jets with $p_T > 25$ GeV and $|\eta| < 2.5$, reconstructed using the anti- k_t algorithm [58,59] with radius parameter $R = 0.4$. Jets from additional collisions in the same bunch crossing are rejected using dedicated tracking and vertex requirements [60]. At least one of the jets must be identified as originating from the decay of a b -hadron (b -jet) using a multivariate tagging algorithm with a 70% efficiency for b -jets and light-quark and gluon mistag rates of 0.5%–2% [61].

Event reconstruction.—Jets are assigned to the top quarks using a χ^2 algorithm that relies on kinematic constraints and the expected values of the top quark and W boson masses [22]. The invariant mass $m_{t\bar{t}}^{\text{reco}}$ of the candidate $t\bar{t}$ pair is reconstructed from the four selected jets, the lepton, and the E_T^{miss} vector. The experimental resolution for the $t\bar{t}$ invariant mass is 8% for $m_{A/H} = 500$ GeV. Events in the $e + \text{jets}$ and $\mu + \text{jets}$ channels are classified into three categories, based on whether a b -tagged jet was assigned to either the hadronically or the semileptonically decaying top quark, or to both of them. Each category defines a signal region; hence six orthogonal signal regions are used in the statistical analysis.

Systematic uncertainties.—The impact of the systematic uncertainties on both the normalization and the shape of the $m_{t\bar{t}}^{\text{reco}}$ distributions is taken into account. The average

TABLE I. Average impact of the dominant uncertainties on the estimated yields for the total background and for a pseudoscalar A with $m_A = 500$ GeV and $\tan\beta = 0.68$ in percent of the nominal value for all signal regions combined. Only uncertainties with a yield impact $> 0.5\%$ are shown. Dots (\dots) indicate that an uncertainty is not applicable to a sample.

Systematic uncertainties [%]	Total background	S	$S + I$
Luminosity [62]	1.7	1.9	1.9
PDF	2.5	2.1	12
$t\bar{t}$ initial-/final-state radiation	3.2	\dots	\dots
$t\bar{t}$ parton shower + fragmentation	4.9	\dots	\dots
$t\bar{t}$ normalization	5.7	\dots	\dots
$t\bar{t}$ event generator	0.5	\dots	\dots
Top quark mass	0.5	2.2	13
Jet energy scale	6.4	4.9	9.3
Jet energy resolution	1.3	1.6	1.7
b -tagging: b -jet efficiency	1.5	1.3	1.1
b -tagging: c -jet efficiency	0.2	0.2	0.8
Electron efficiency	0.3	0.4	0.7
Muon efficiency	0.9	1.0	1.0
Signal MC scales	\dots	7.3	7.3
Reweighting	\dots	\dots	5.0
MC statistical uncertainty	0.5	2.4	11
Total uncertainty	11	10	25

impact of the dominant uncertainties on the event yields is summarized in Table I.

The experimental uncertainties with the largest impact on the event yields and the shape of the $m_{t\bar{t}}^{\text{reco}}$ distributions are those related to the jet energy scale (JES) and the jet energy resolution (JER) [63,64], followed by uncertainties on the b -tagging efficiency and misidentification rates [61]. The uncertainties related to leptons include those in the reconstruction and isolation efficiency, the single-lepton triggers, and the energy scale and resolution [54,55].

The uncertainty of 6.5% in the NNLO + NNLL cross section for SM $t\bar{t}$ production is the dominant uncertainty in the total background normalization [22]. Modeling uncertainties affecting the shape of the $m_{t\bar{t}}^{\text{reco}}$ distribution for the SM $t\bar{t}$ background are also taken into account. These

TABLE II. Number of events observed in data and expected number of background events after the event selection, before the profile-likelihood fit to the full data set. The uncertainty in the background yields is derived by summing all uncertainties in quadrature. The “other bkg.” component comprises single top quark, $t\bar{t} + W/Z$, $Z + \text{jets}$, diboson, and multijet production.

Type	$e + \text{jets}$	$\mu + \text{jets}$
$t\bar{t}$	$95\,000 \pm 11\,000$	$93\,000 \pm 11\,000$
$W + \text{jets}$	6600 ± 2100	7200 ± 2300
Other bkg.	$11\,200 \pm 1400$	6100 ± 600
Total	$112\,800 \pm 13\,000$	$106\,300 \pm 12\,000$
Data	115 785	110 218

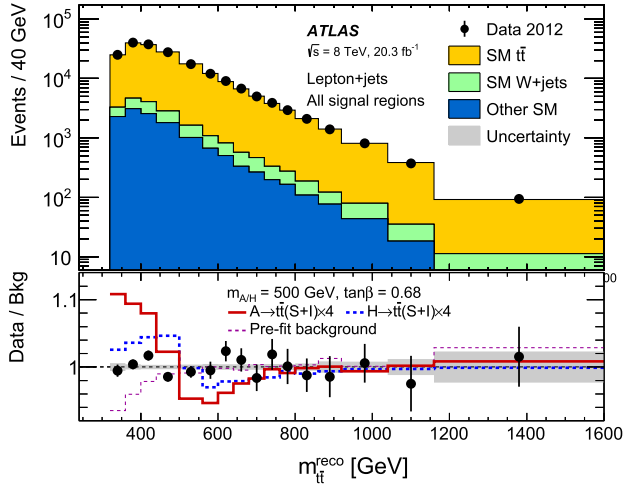


FIG. 2. Distribution of $m_{t\bar{t}}^{\text{reco}}$ for the data and the expected background after the profile-likelihood fit under the background-only hypothesis for all signal regions combined. The lines in the bottom panel show the individual $S + I$ distributions (scaled by a factor of 4) for a pseudoscalar A (solid line) and scalar H (bold dashed line) with $m_{A/H} = 500 \text{ GeV}$ and $\tan\beta = 0.68$ relative to the total background.

include uncertainties related to the choice of NLO event generator, the modeling of the parton shower and fragmentation, the modeling of gluon initial- and final-state radiation, and the value of the top quark mass m_t . Other sources of uncertainty related to the various background components are described in Ref. [22].

The largest uncertainty in the modeling of the $S + I$ and S components is related to the $\pm 1.0 \text{ GeV}$ uncertainty of the value $m_t = 172.5 \text{ GeV}$ [65]. Uncertainties related to the choice of PDF set and renormalization and factorization scales are also considered. The latter is estimated by varying the scales by factors of 0.5 and 2.0, which yields a constant $\pm 7.3\%$ variation across the $m_{t\bar{t}}^{\text{reco}}$ spectrum. An asymmetric variation, for which the bins at the low and high ends of the $m_{t\bar{t}}^{\text{reco}}$ spectrum are taken as anticorrelated [66] is also considered to estimate the impact of the scale variations on the shape of the $m_{t\bar{t}}^{\text{reco}}$ spectrum. For the

$S + I$ samples, an additional constant $\pm 5\%$ uncertainty is included to cover the difference between reweighted and generated distributions.

Results.—A breakdown of the observed and expected event yields in the $e + \text{jets}$ and $\mu + \text{jets}$ channels and their total uncertainties is shown in Table II. Good agreement is found between the observed number of events in data and the expected total number of background events.

The exclusion limits are derived separately for each signal hypothesis from a profile-likelihood fit [67] of the expected $m_{t\bar{t}}^{\text{reco}}$ distributions to the observed ones simultaneously in all signal regions, taking the statistical and systematic uncertainties into account as nuisance parameters [22]. Only bins with $m_{t\bar{t}}^{\text{reco}} > 320 \text{ GeV}$ are considered to avoid threshold effects not well described by the simulation. The shape of the binned $m_{t\bar{t}}^{\text{reco}}$ distributions is parametrized in terms of the signal strength μ [26,27]:

$$\mu S + \sqrt{\mu} I + B = (\mu - \sqrt{\mu}) S + \sqrt{\mu} (S + I) + B. \quad (1)$$

The fitted variable is $\sqrt{\mu}$ and the case $\mu = 1$ ($\mu = 0$) corresponds to the type-II 2HDM in the alignment limit (the background-only hypothesis). This approach relies on the assumption that, for a given signal hypothesis, the shape of the $t\bar{t}$ invariant mass distributions for S and $S + I$ in Eq. (1) does not change with μ . The terms S and $S + I$ on the right-hand side of Eq. (1) correspond to the $m_{t\bar{t}}^{\text{reco}}$ distributions obtained from the S and $S + I$ samples, respectively, while B stands for the expected $m_{t\bar{t}}^{\text{reco}}$ distribution of the total background.

The level of agreement between the observed and expected mass spectra is quantified in a fit under the background-only hypothesis in which only the nuisance parameters are allowed to vary. The observed $m_{t\bar{t}}^{\text{reco}}$ spectra are compatible with the postfit expected spectra within the (constrained) uncertainty bands (Fig. 2).

The upper limits on μ at 95% confidence level (C.L.) are obtained with the C.L._s method [68] for a number of $(m_{A/H}, \tan\beta)$ values. The upper limits at intermediate points are obtained from a linear interpolation among

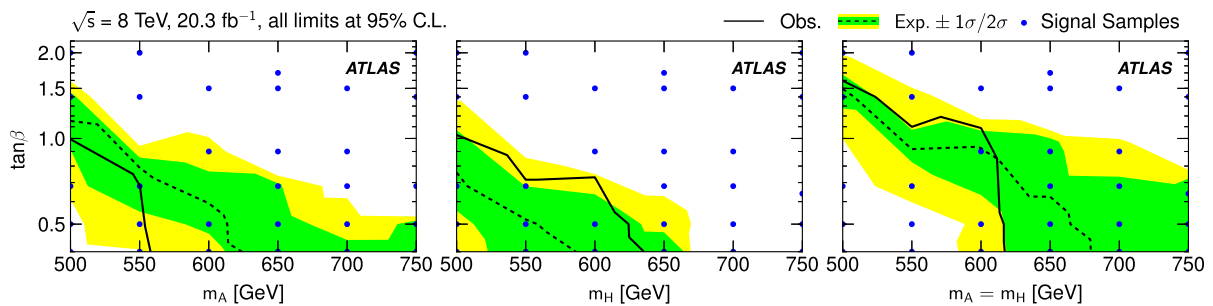


FIG. 3. The 95% C.L. observed and expected exclusion regions for the type-II 2HDM ($\mu = 1$) considering only a pseudoscalar A (left), only a scalar H (middle), and the mass-degenerate scenario $m_A = m_H$ (right). Blue points indicate parameter values at which signal samples are produced.

TABLE III. The 95% C.L. observed (obs.) and expected (exp.) exclusion limits on $\tan\beta$ for a type-II 2HDM in the alignment limit considering only a pseudoscalar A (left), only a scalar H (middle), and the mass-degenerate scenario $m_A = m_H$ (right). Dots (\dots) indicate that no value of $\tan\beta \geq 0.4$ is excluded.

Mass [GeV]	$\tan\beta$:	m_A		m_H		$m_A = m_H$	
		obs.	exp.	obs.	exp.	obs.	exp.
500		<1.00	<1.16	<1.00	<0.77	<1.55	<1.50
550		<0.69	<0.79	<0.72	<0.52	<1.10	<0.92
600		\dots	<0.59	<0.73	\dots	<1.09	<0.93
650		\dots	\dots	\dots	\dots	\dots	<0.62

the three closest points. In Fig. 3, the observed and expected exclusion regions for the type-II 2HDM ($\mu = 1$) are shown for the three scenarios discussed in the Introduction. The excluded values of $\tan\beta$ for the different mass hypotheses are listed in Table III.

Conclusion.—In conclusion, the search for massive pseudoscalar and scalar resonances decaying to $t\bar{t}$ in 20.3 fb^{-1} of pp collisions at 8 TeV recorded by the ATLAS experiment yields no statistically significant deviations from the SM prediction. The results are interpreted in a type-II 2HDM in the alignment limit, and upper limits are set on the signal strength μ at 95% C.L. in the $m_{A/H}$ versus $\tan\beta$ plane. Unlike previous searches for $t\bar{t}$ resonances, this analysis takes into account interference effects between the signal process and the background from SM $t\bar{t}$ production. It tightens significantly the previously published constraints on the 2HDM parameter space in the low $\tan\beta$ and high mass ($m_{A/H} > 500 \text{ GeV}$) region.

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Cervelli,¹⁸ S. A. Cetin,^{20d} A. Chafaq,^{137a} D. Chakraborty,¹¹⁰ S. K. Chan,⁵⁹ W. S. Chan,¹⁰⁹ Y. L. Chan,^{62a} P. Chang,¹⁶⁹ J. D. Chapman,³⁰ D. G. Charlton,¹⁹ C. C. Chau,³¹ C. A. Chavez Barajas,¹⁵¹ S. Che,¹¹³ S. Cheatham,^{167a,167c} A. Chegwidan,⁹³ S. Chekanov,⁶ S. V. Chekulaev,^{163a} G. A. Chelkov,^{68,m} M. A. Chelstowska,³² C. Chen,⁶⁷ H. Chen,²⁷ J. Chen,^{36a} S. Chen,^{35b} S. Chen,¹⁵⁷ X. Chen,^{35c,n} Y. Chen,⁷⁰ Y. Chen,⁴⁵ H. C. Cheng,⁹² H. J. Cheng,^{35a} A. Cheplakov,⁶⁸ E. Cheremushkina,¹³² R. Cherkaoui El Moursli,^{137e} E. Cheu,⁷ K. Cheung,⁶³ L. Chevalier,¹³⁸ V. Chiarella,⁵⁰ G. Chiarelli,^{126a,126b} G. Chiodini,^{76a} A. S. Chisholm,³² A. Chitan,^{28b} Y. H. Chiu,¹⁷² M. V. Chizhov,⁶⁸ K. Choi,⁶⁴ A. R. Chomont,³⁷ S. Chouridou,¹⁵⁶ Y. S. Chow,^{62a} V. Christodoulou,⁸¹ M. C. Chu,^{62a} J. Chudoba,¹²⁹ A. J. Chuinard,⁹⁰ J. J. Chwastowski,⁴² L. Chytka,¹¹⁷ A. K. Ciftci,^{4a} D. Cinca,⁴⁶ V. Cindro,⁷⁸ I. A. Cioara,²³ C. Ciocca,^{22a,22b} A. Ciocio,¹⁶ F. Ciroto,^{106a,106b} Z. H. Citron,¹⁷⁵ M. Citterio,^{94a} M. Ciubancan,^{28b} A. Clark,⁵² B. L. Clark,⁵⁹ M. R. Clark,³⁸ P. J. Clark,⁴⁹

R. N. Clarke,¹⁶ C. Clement,^{148a,148b} Y. Coadou,⁸⁸ M. Cobal,^{167a,167c} A. Coccaro,⁵² J. Cochran,⁶⁷ L. Colasurdo,¹⁰⁸ B. Cole,³⁸ A. P. Colijn,¹⁰⁹ J. Collot,⁵⁸ T. Colombo,¹⁶⁶ P. Conde Muiño,^{128a,128b} E. Coniavitis,⁵¹ S. H. Connell,^{147b} I. A. Connelly,⁸⁷ S. Constantinescu,^{28b} G. Conti,³² F. Conventi,^{106a,o} M. Cooke,¹⁶ A. M. Cooper-Sarkar,¹²² F. Cormier,¹⁷¹ K. J. R. Cormier,¹⁶¹ M. Corradi,^{134a,134b} F. Corriveau,^{90,p} A. Cortes-Gonzalez,³² G. Cortiana,¹⁰³ G. Costa,^{94a} M. J. Costa,¹⁷⁰ D. Costanzo,¹⁴¹ G. Cottin,³⁰ G. Cowan,⁸⁰ B. E. Cox,⁸⁷ K. Cranmer,¹¹² S. J. Crawley,⁵⁶ R. A. Creager,¹²⁴ G. Cree,³¹ S. Crépé-Renaudin,⁵⁸ F. Crescioli,⁸³ W. A. Cribbs,^{148a,148b} M. Cristinziani,²³ V. Croft,¹⁰⁸ G. Crosetti,^{40a,40b} A. Cueto,⁸⁵ T. Cuhadar Donszelmann,¹⁴¹ A. R. Cukierman,¹⁴⁵ J. Cummings,¹⁷⁹ M. Curatolo,⁵⁰ J. Cúth,⁸⁶ S. Czekierda,⁴² P. Czodrowski,³² G. D'amen,^{22a,22b} S. D'Auria,⁵⁶ L. D'eraimo,⁸³ M. D'Onofrio,⁷⁷ M. J. Da Cunha Sargedas De Sousa,^{128a,128b} C. Da Via,⁸⁷ W. Dabrowski,^{41a} T. Dado,^{146a} T. Dai,⁹² O. Dale,¹⁵ F. Dallaire,⁹⁷ C. Dallapiccola,⁸⁹ M. Dam,³⁹ J. R. Dandoy,¹²⁴ M. F. Daneri,²⁹ N. P. Dang,¹⁷⁶ A. C. Daniells,¹⁹ N. S. Dann,⁸⁷ M. Danninger,¹⁷¹ M. Dano Hoffmann,¹³⁸ V. Dao,¹⁵⁰ G. Darbo,^{53a} S. Darmora,⁸ J. Dassoulas,³ A. Dattagupta,¹¹⁸ T. Daubney,⁴⁵ W. Davey,²³ C. David,⁴⁵ T. Davidek,¹³¹ D. R. Davis,⁴⁸ P. Davison,⁸¹ E. Dawe,⁹¹ I. Dawson,¹⁴¹ K. De,⁸ R. de Asmundis,^{106a} A. De Benedetti,¹¹⁵ S. De Castro,^{22a,22b} S. De Cecco,⁸³ N. De Groot,¹⁰⁸ P. de Jong,¹⁰⁹ H. De la Torre,⁹³ F. De Lorenzi,⁶⁷ A. De Maria,⁵⁷ D. De Pedis,^{134a} A. De Salvo,^{134a} U. De Sanctis,^{135a,135b} A. De Santo,¹⁵¹ K. De Vasconcelos Corga,⁸⁸ J. B. De Vivie De Regie,¹¹⁹ W. J. Dearnaley,⁷⁵ R. Debbe,²⁷ C. Debenedetti,¹³⁹ D. V. Dedovich,⁶⁸ N. Dehghanian,³ I. Deigaard,¹⁰⁹ M. Del Gaudio,^{40a,40b} J. Del Peso,⁸⁵ D. Delgove,¹¹⁹ F. Deliot,¹³⁸ C. M. Delitzsch,⁷ A. Dell'Acqua,³² L. Dell'Asta,²⁴ M. Dell'Orso,^{126a,126b} M. Della Pietra,^{106a,106b} D. della Volpe,⁵² M. Delmastro,⁵ C. Delporte,¹¹⁹ P. A. Delsart,⁵⁸ D. A. DeMarco,¹⁶¹ S. Demers,¹⁷⁹ M. Demichev,⁶⁸ A. Demilly,⁸³ S. P. Denisov,¹³² D. Denysiuk,¹³⁸ D. Derendarz,⁴² J. E. Derkaoui,^{137d} F. Derue,⁸³ P. Dervan,⁷⁷ K. Desch,²³ C. Deterre,⁴⁵ K. Dette,⁴⁶ M. R. Devesa,²⁹ P. O. Deviveiros,³² A. Dewhurst,¹³³ S. Dhaliwal,²⁵ F. A. Di Bello,⁵² A. Di Ciaccio,^{135a,135b} L. Di Ciaccio,⁵ W. K. Di Clemente,¹²⁴ C. Di Donato,^{106a,106b} A. Di Girolamo,³² B. Di Girolamo,³² B. Di Micco,^{136a,136b} R. Di Nardo,³² K. F. Di Petrillo,⁵⁹ A. Di Simone,⁵¹ R. Di Sipio,¹⁶¹ D. Di Valentino,³¹ C. Diaconu,⁸⁸ M. Diamond,¹⁶¹ F. A. Dias,³⁹ M. A. Diaz,^{34a} E. B. Diehl,⁹² J. Dietrich,¹⁷ S. Díez Cornell,⁴⁵ A. Dimitrievska,¹⁴ J. Dingfelder,²³ P. Dita,^{28b} S. Dita,^{28b} F. Dittus,³² F. Djama,⁸⁸ T. Djobava,^{54b} J. I. Djuvsland,^{60a} M. A. B. do Vale,^{26c} D. Dobos,³² M. Dobre,^{28b} C. Doglioni,⁸⁴ J. Dolejsi,¹³¹ Z. Dolezal,¹³¹ M. Donadelli,^{26d} S. Donati,^{126a,126b} P. Dondero,^{123a,123b} J. Donini,³⁷ J. Dopke,¹³³ A. Doria,^{106a} M. T. Dova,⁷⁴ A. T. Doyle,⁵⁶ E. Drechsler,⁵⁷ M. Dris,¹⁰ Y. Du,^{36b} J. Duarte-Camperderros,¹⁵⁵ A. Dubreuil,⁵² E. Duchovni,¹⁷⁵ G. Duckeck,¹⁰² A. Ducourthial,⁸³ O. A. Ducu,^{97,q} D. Duda,¹⁰⁹ A. Dudarev,³² A. Chr. Dudder,⁸⁶ E. M. Duffield,¹⁶ L. Dufлот,¹¹⁹ M. Dührssen,³² M. Dumancic,¹⁷⁵ A. E. Dumitriu,^{28b} A. K. Duncan,⁵⁶ M. Dunford,^{60a} H. Duran Yildiz,^{4a} M. Düren,⁵⁵ A. Durglishvili,^{54b} D. Duschinger,⁴⁷ B. Dutta,⁴⁵ D. Duvnjak,¹ M. Dyndal,⁴⁵ B. S. Dziedzic,⁴² C. Eckardt,⁴⁵ K. M. Ecker,¹⁰³ R. C. Edgar,⁹² T. Eifert,³² G. Eigen,¹⁵ K. Einsweiler,¹⁶ T. Ekelof,¹⁶⁸ M. El Kacimi,^{137c} R. El Kosseifi,⁸⁸ V. Ellajosyula,⁸⁸ M. Ellert,¹⁶⁸ S. Elles,⁵ F. Ellinghaus,¹⁷⁸ A. A. Elliot,¹⁷² N. Ellis,³² J. Elmsheuser,²⁷ M. Elsing,³² D. Emeliyanov,¹³³ Y. Enari,¹⁵⁷ O. C. Endner,⁸⁶ J. S. Ennis,¹⁷³ J. Erdmann,⁴⁶ A. Ereditato,¹⁸ M. Ernst,²⁷ S. Errede,¹⁶⁹ M. Escalier,¹¹⁹ C. Escobar,¹⁷⁰ B. Esposito,⁵⁰ O. Estrada Pastor,¹⁷⁰ A. I. Etienvre,¹³⁸ E. Etzion,¹⁵⁵ H. Evans,⁶⁴ A. Ezhilov,¹²⁵ M. Ezzi,^{137e} F. Fabbri,^{22a,22b} L. Fabbri,^{22a,22b} V. Fabiani,¹⁰⁸ G. Facini,⁸¹ R. M. Fakhruddinov,¹³² S. Falciano,^{134a} R. J. Falla,⁸¹ J. Faltova,³² Y. Fang,^{35a} M. Fanti,^{94a,94b} A. Farbin,⁸ A. Farilla,^{136a} C. Farina,¹²⁷ E. M. Farina,^{123a,123b} T. Farooque,⁹³ S. Farrell,¹⁶ S. M. Farrington,¹⁷³ P. Farthouat,³² F. Fassi,^{137e} P. Fassnacht,³² D. Fassouliotis,⁹ M. Faucci Giannelli,⁴⁹ A. Favareto,^{53a,53b} W. J. Fawcett,¹²² L. Fayard,¹¹⁹ O. L. Fedin,^{125,r} W. Fedorko,¹⁷¹ S. Feigl,¹²¹ L. Felgioni,⁸⁸ C. Feng,^{36b} E. J. Feng,³² H. Feng,⁹² M. J. Fenton,⁵⁶ A. B. Fenyuk,¹³² L. Feremenga,⁸ P. Fernandez Martinez,¹⁷⁰ S. Fernandez Perez,¹³ J. Ferrando,⁴⁵ A. Ferrari,¹⁶⁸ P. Ferrari,¹⁰⁹ R. Ferrari,^{123a} D. E. Ferreira de Lima,^{60b} A. Ferrer,¹⁷⁰ D. Ferrere,⁵² C. Ferretti,⁹² F. Fiedler,⁸⁶ A. Filipčič,⁷⁸ M. Filipuzzi,⁴⁵ F. Filthaut,¹⁰⁸ M. Fincke-Keeler,¹⁷² K. D. Finelli,¹⁵² M. C. N. Fiolhais,^{128a,128c,s} L. Fiorini,¹⁷⁰ A. Fischer,² C. Fischer,¹³ J. Fischer,¹⁷⁸ W. C. Fisher,⁹³ N. Flaschel,⁴⁵ I. Fleck,¹⁴³ P. Fleischmann,⁹² R. R. M. Fletcher,¹²⁴ T. Flick,¹⁷⁸ B. M. Flierl,¹⁰² L. R. Flores Castillo,^{62a} M. J. Flowerdew,¹⁰³ G. T. Forcolin,⁸⁷ A. Formica,¹³⁸ F. A. Förster,¹³ A. Forti,⁸⁷ A. G. Foster,¹⁹ D. Fournier,¹¹⁹ H. Fox,⁷⁵ S. Fracchia,¹⁴¹ P. Francavilla,⁸³ M. Franchini,^{22a,22b} S. Franchino,^{60a} D. Francis,³² L. Franconi,¹²¹ M. Franklin,⁵⁹ M. Frate,¹⁶⁶ M. Fraternali,^{123a,123b} D. Freeborn,⁸¹ S. M. Fressard-Batraneanu,³² B. Freund,⁹⁷ D. Froidevaux,³² J. A. Frost,¹²² C. Fukunaga,¹⁵⁸ T. Fusayasu,¹⁰⁴ J. Fuster,¹⁷⁰ C. Gabaldon,⁵⁸ O. Gabizon,¹⁵⁴ A. Gabrielli,^{22a,22b} A. Gabrielli,¹⁶ G. P. Gach,^{41a} S. Gadatsch,³² S. Gadomski,⁸⁰ G. Gagliardi,^{53a,53b} L. G. Gagnon,⁹⁷ C. Galea,¹⁰⁸ B. Galhardo,^{128a,128c} E. J. Gallas,¹²² B. J. Gallop,¹³³ P. Gallus,¹³⁰ G. Galster,³⁹ K. K. Gan,¹¹³ S. Ganguly,³⁷ Y. Gao,⁷⁷ Y. S. Gao,^{145,h} F. M. Garay Walls,⁴⁹ C. García,¹⁷⁰ J. E. García Navarro,¹⁷⁰ J. A. García Pascual,^{35a} M. Garcia-Sciveres,¹⁶ R. W. Gardner,³³ N. Garelli,¹⁴⁵ V. Garonne,¹²¹ A. Gascon Bravo,⁴⁵ K. Gasnikova,⁴⁵ C. Gatti,⁵⁰ A. Gaudiello,^{53a,53b} G. Gaudio,^{123a}

I. L. Gavrilenko,⁹⁸ C. Gay,¹⁷¹ G. Gaycken,²³ E. N. Gazis,¹⁰ C. N. P. Gee,¹³³ J. Geisen,⁵⁷ M. Geisen,⁸⁶ M. P. Geisler,^{60a} K. Gellerstedt,^{148a,148b} C. Gemme,^{53a} M. H. Genest,⁵⁸ C. Geng,⁹² S. Gentile,^{134a,134b} C. Gentsos,¹⁵⁶ S. George,⁸⁰ D. Gerbaudo,¹³ A. Gershon,¹⁵⁵ G. Geßner,⁴⁶ S. Ghasemi,¹⁴³ M. Ghneimat,²³ B. Giacobbe,^{22a} S. Giagu,^{134a,134b} N. Giangiacomi,^{22a,22b} P. Giannetti,^{126a,126b} S. M. Gibson,⁸⁰ M. Gignac,¹⁷¹ M. Gilchriese,¹⁶ D. Gillberg,³¹ G. Gilles,¹⁷⁸ D. M. Gingrich,^{3,e} N. Giokaris,^{9,a} M. P. Giordani,^{167a,167c} F. M. Giorgi,^{22a} P. F. Giraud,¹³⁸ P. Giromini,⁵⁹ G. Giugliarelli,^{167a,167c} D. Giugni,^{94a} F. Giuli,¹²² C. Giuliani,¹⁰³ M. Giulini,^{60b} B. K. Gjelsten,¹²¹ S. Gkaitatzis,¹⁵⁶ I. Gkialas,^{9,t} E. L. Gkougkousis,¹³⁹ P. Gkoutoumis,¹⁰ L. K. Gladilin,¹⁰¹ C. Glasman,⁸⁵ J. Glatzer,¹³ P. C. F. Glaysheer,⁴⁵ A. Glazov,⁴⁵ M. Goblirsch-Kolb,²⁵ J. Godlewski,⁴² S. Goldfarb,⁹¹ T. Golling,⁵² D. Golubkov,¹³² A. Gomes,^{128a,128b,128d} R. Gonçalves,^{128a} R. Goncalves Gama,^{26a} J. Goncalves Pinto Firmino Da Costa,¹³⁸ G. Gonella,⁵¹ L. Gonella,¹⁹ A. Gongadze,⁶⁸ S. González de la Hoz,¹⁷⁰ S. Gonzalez-Sevilla,⁵² L. Goossens,³² P. A. Gorbounov,⁹⁹ H. A. Gordon,²⁷ I. Gorelov,¹⁰⁷ B. Gorini,³² E. Gorini,^{76a,76b} A. Gorišek,⁷⁸ A. T. Goshaw,⁴⁸ C. Gössling,⁴⁶ M. I. Gostkin,⁶⁸ C. A. Gottardo,²³ C. R. Goudet,¹¹⁹ D. Goujdami,^{137c} A. G. Goussiou,¹⁴⁰ N. Govender,^{147b,u} E. Gozani,¹⁵⁴ L. Graber,⁵⁷ I. Grabowska-Bold,^{41a} P. O. J. Gradin,¹⁶⁸ J. Gramling,¹⁶⁶ E. Gramstad,¹²¹ S. Grancagnolo,¹⁷ V. Gratchev,¹²⁵ P. M. Gravila,^{28f} C. Gray,⁵⁶ H. M. Gray,¹⁶ Z. D. Greenwood,^{82,v} C. Greife,²³ K. Gregersen,⁸¹ I. M. Gregor,⁴⁵ P. Grenier,¹⁴⁵ K. Grevtsov,⁵ J. Griffiths,⁸ A. A. Grillo,¹³⁹ K. Grimm,⁷⁵ S. Grinstein,^{13,w} Ph. Gris,³⁷ J.-F. Grivaz,¹¹⁹ S. Groh,⁸⁶ E. Gross,¹⁷⁵ J. Grosse-Knetter,⁵⁷ G. C. Grossi,⁸² Z. J. Grout,⁸¹ A. Grummer,¹⁰⁷ L. Guan,⁹² W. Guan,¹⁷⁶ J. Guenther,⁶⁵ F. Guescini,^{163a} D. Guest,¹⁶⁶ O. Gueta,¹⁵⁵ B. Gui,¹¹³ E. Guido,^{53a,53b} T. Guillemin,⁵ S. Guindon,² U. Gul,⁵⁶ C. Gumpert,³² J. Guo,^{36c} W. Guo,⁹² Y. Guo,^{36a} R. Gupta,⁴³ S. Gupta,¹²² G. Gustavino,¹¹⁵ P. Gutierrez,¹¹⁵ N. G. Gutierrez Ortiz,⁸¹ C. Gutsche,⁸¹ C. Guyot,¹³⁸ M. P. Guzik,^{41a} C. Gwenlan,¹²² C. B. Gwilliam,⁷⁷ A. Haas,¹¹² C. Haber,¹⁶ H. K. Hadavand,⁸ N. Haddad,^{137e} A. Hadeef,⁸⁸ S. Hageböck,²³ M. Hagihara,¹⁶⁴ H. Hakobyan,^{180,a} M. Haleem,⁴⁵ J. Haley,¹¹⁶ G. Halladjian,⁹³ G. D. Hallewell,⁸⁸ K. Hamacher,¹⁷⁸ P. Hamal,¹¹⁷ K. Hamano,¹⁷² A. Hamilton,^{147a} G. N. Hamity,¹⁴¹ P. G. Hamnett,⁴⁵ L. Han,^{36a} S. Han,^{35a} K. Hanagaki,^{69,x} K. Hanawa,¹⁵⁷ M. Hance,¹³⁹ B. Haney,¹²⁴ P. Hanke,^{60a} J. B. Hansen,³⁹ J. D. Hansen,³⁹ M. C. Hansen,²³ P. H. Hansen,³⁹ K. Hara,¹⁶⁴ A. S. Hard,¹⁷⁶ T. Harenberg,¹⁷⁸ F. Hariri,¹¹⁹ S. Harkusha,⁹⁵ R. D. Harrington,⁴⁹ P. F. Harrison,¹⁷³ N. M. Hartmann,¹⁰² M. Hasegawa,⁷⁰ Y. Hasegawa,¹⁴² A. Hasib,⁴⁹ S. Hassani,¹³⁸ S. Haug,¹⁸ R. Hauser,⁹³ L. Hauswald,⁴⁷ L. B. Havener,³⁸ M. Havranek,¹³⁰ C. M. Hawkes,¹⁹ R. J. Hawkings,³² D. Hayakawa,¹⁵⁹ D. Hayden,⁹³ C. P. Hays,¹²² J. M. Hays,⁷⁹ H. S. Hayward,⁷⁷ S. J. Haywood,¹³³ S. J. Head,¹⁹ T. Heck,⁸⁶ V. Hedberg,⁸⁴ L. Heelan,⁸ S. Heer,²³ K. K. Heidegger,⁵¹ S. Heim,⁴⁵ T. Heim,¹⁶ B. Heinemann,^{45,y} J. J. Heinrich,¹⁰² L. Heinrich,¹¹² C. Heinz,⁵⁵ J. Hejbal,¹²⁹ L. Helary,³² A. Held,¹⁷¹ S. Hellman,^{148a,148b} C. Hensens,³² R. C. W. Henderson,⁷⁵ Y. Heng,¹⁷⁶ S. Henkelmann,¹⁷¹ A. M. Henriques Correia,³² S. Henrot-Versille,¹¹⁹ G. H. Herbert,¹⁷ H. Herde,²⁵ V. Herget,¹⁷⁷ Y. Hernández Jiménez,^{147c} H. Herr,⁸⁶ G. Herten,⁵¹ R. Hertenberger,¹⁰² L. Hervas,³² T. C. Herwig,¹²⁴ G. G. Hesketh,⁸¹ N. P. Hessey,^{163a} J. W. Hetherly,⁴³ S. Higashino,⁶⁹ E. Higón-Rodriguez,¹⁷⁰ K. Hildebrand,³³ E. Hill,¹⁷² J. C. Hill,³⁰ K. H. Hiller,⁴⁵ S. J. Hillier,¹⁹ M. Hils,⁴⁷ I. Hinchliffe,¹⁶ M. Hirose,⁵¹ D. Hirschbuehl,¹⁷⁸ B. Hiti,⁷⁸ O. Hladik,¹²⁹ X. Hoad,⁴⁹ J. Hobbs,¹⁵⁰ N. Hod,^{163a} M. C. Hodgkinson,¹⁴¹ P. Hodgson,¹⁴¹ A. Hoecker,³² M. R. Hoferkamp,¹⁰⁷ F. Hoenic,¹⁰² D. Hohn,²³ T. R. Holmes,³³ M. Homann,⁴⁶ S. Honda,¹⁶⁴ T. Honda,⁶⁹ T. M. Hong,¹²⁷ B. H. Hooberman,¹⁶⁹ W. H. Hopkins,¹¹⁸ Y. Horii,¹⁰⁵ A. J. Horton,¹⁴⁴ J.-Y. Hostachy,⁵⁸ S. Hou,¹⁵³ A. Hoummada,^{137a} J. Howarth,⁸⁷ J. Hoya,⁷⁴ M. Hrabovsky,¹¹⁷ J. Hrdinka,³² I. Hristova,¹⁷ J. Hrivnac,¹¹⁹ T. Hryn'ova,⁵ A. Hrynevich,⁹⁶ P. J. Hsu,⁶³ S.-C. Hsu,¹⁴⁰ Q. Hu,^{36a} S. Hu,^{36c} Y. Huang,^{35a} Z. Hubacek,¹³⁰ F. Hubaut,⁸⁸ F. Huegging,²³ T. B. Huffman,¹²² E. W. Hughes,³⁸ G. Hughes,⁷⁵ M. Huhtinen,³² P. Huo,¹⁵⁰ N. Huseynov,^{68,c} J. Huston,⁹³ J. Huth,⁵⁹ G. Iacobucci,⁵² G. Iakovidis,²⁷ I. Ibragimov,¹⁴³ L. Iconomidou-Fayard,¹¹⁹ Z. Idrissi,^{137e} P. Iengo,³² O. Igonkina,^{109,z} T. Iizawa,¹⁷⁴ Y. Ikegami,⁶⁹ M. Ikeno,⁶⁹ Y. Ilchenko,^{11,aa} D. Iliadis,¹⁵⁶ N. Ilic,¹⁴⁵ G. Introzzi,^{123a,123b} P. Ioannou,^{9,a} M. Iodice,^{136a} K. Iordanidou,³⁸ V. Ippolito,⁵⁹ M. F. Isacson,¹⁶⁸ N. Ishijima,¹²⁰ M. Ishino,¹⁵⁷ M. Ishitsuka,¹⁵⁹ C. Issever,¹²² S. Istin,^{20a} F. Ito,¹⁶⁴ J. M. Iturbe Ponce,^{62a} R. Iuppa,^{162a,162b} H. Iwasaki,⁶⁹ J. M. Izen,⁴⁴ V. Izzo,^{106a} S. Jabbar,³ P. Jackson,¹ R. M. Jacobs,²³ V. Jain,² K. B. Jakobi,⁸⁶ K. Jakobs,⁵¹ S. Jakobsen,⁶⁵ T. Jakoubek,¹²⁹ D. O. Jamin,¹¹⁶ D. K. Jana,⁸² R. Jansky,⁵² J. Janssen,²³ M. Janus,⁵⁷ P. A. Janus,^{41a} G. Jarlskog,⁸⁴ N. Javadov,^{68,c} T. Javůrek,⁵¹ M. Javurkova,⁵¹ F. Jeanneau,¹³⁸ L. Jeanty,¹⁶ J. Jejelava,^{54a,bb} A. Jelinskas,¹⁷³ P. Jenni,^{51,cc} C. Jeske,¹⁷³ S. Jézéquel,⁵ H. Ji,¹⁷⁶ J. Jia,¹⁵⁰ H. Jiang,⁶⁷ Y. Jiang,^{36a} Z. Jiang,¹⁴⁵ S. Jiggins,⁸¹ J. Jimenez Pena,¹⁷⁰ S. Jin,^{35a} A. Jinaru,^{28b} O. Jinnouchi,¹⁵⁹ H. Jivan,^{147c} P. Johansson,¹⁴¹ K. A. Johns,⁷ C. A. Johnson,⁶⁴ W. J. Johnson,¹⁴⁰ K. Jon-And,^{148a,148b} R. W. L. Jones,⁷⁵ S. D. Jones,¹⁵¹ S. Jones,⁷ T. J. Jones,⁷⁷ J. Jongmanns,^{60a} P. M. Jorge,^{128a,128b} J. Jovicevic,^{163a} X. Ju,¹⁷⁶ A. Juste Rozas,^{13,w} M. K. Köhler,¹⁷⁵ A. Kaczmarska,⁴² M. Kado,¹¹⁹ H. Kagan,¹¹³ M. Kagan,¹⁴⁵ S. J. Kahn,⁸⁸ T. Kaji,¹⁷⁴ E. Kajomovitz,⁴⁸ C. W. Kalderon,⁸⁴ A. Kaluza,⁸⁶ S. Kama,⁴³ A. Kamenshchikov,¹³² N. Kanaya,¹⁵⁷ L. Kanjir,⁷⁸ V. A. Kantserov,¹⁰⁰ J. Kanzaki,⁶⁹ B. Kaplan,¹¹²

L. S. Kaplan,¹⁷⁶ D. Kar,^{147c} K. Karakostas,¹⁰ N. Karastathis,¹⁰ M. J. Kareem,⁵⁷ E. Karentzos,¹⁰ S. N. Karpov,⁶⁸
 Z. M. Karpova,⁶⁸ K. Karthik,¹¹² V. Kartvelishvili,⁷⁵ A. N. Karyukhin,¹³² K. Kasahara,¹⁶⁴ L. Kashif,¹⁷⁶ R. D. Kass,¹¹³
 A. Kastanas,¹⁴⁹ Y. Kataoka,¹⁵⁷ C. Kato,¹⁵⁷ A. Katre,⁵² J. Katzy,⁴⁵ K. Kawade,⁷⁰ K. Kawagoe,⁷³ T. Kawamoto,¹⁵⁷
 G. Kawamura,⁵⁷ E. F. Kay,⁷⁷ V. F. Kazanin,^{111,d} R. Keeler,¹⁷² R. Kehoe,⁴³ J. S. Keller,³¹ E. Kellermann,⁸⁴ J. J. Kempster,⁸⁰
 J. Kendrick,¹⁹ H. Keoshkerian,¹⁶¹ O. Kepka,¹²⁹ B. P. Kerševan,⁷⁸ S. Kersten,¹⁷⁸ R. A. Keyes,⁹⁰ M. Khader,¹⁶⁹
 F. Khalil-zada,¹² A. Khanov,¹¹⁶ A. G. Kharlamov,^{111,d} T. Kharlamova,^{111,d} A. Khodinov,¹⁶⁰ T. J. Khoo,⁵² V. Khovanskiy,^{99,a}
 E. Khramov,⁶⁸ J. Khubua,^{54b,dd} S. Kido,⁷⁰ C. R. Kilby,⁸⁰ H. Y. Kim,⁸ S. H. Kim,¹⁶⁴ Y. K. Kim,³³ N. Kimura,¹⁵⁶ O. M. Kind,¹⁷
 B. T. King,⁷⁷ D. Kirchmeier,⁴⁷ J. Kirk,¹³³ A. E. Kiryunin,¹⁰³ T. Kishimoto,¹⁵⁷ D. Kisielewska,^{41a} V. Kitali,⁴⁵ K. Kiuchi,¹⁶⁴
 O. Kivernykh,⁵ E. Kladiva,^{146b} T. Klapdor-Kleingrothaus,⁵¹ M. H. Klein,⁹² M. Klein,⁷⁷ U. Klein,⁷⁷ K. Kleinknecht,⁸⁶
 P. Klimek,¹¹⁰ A. Klimentov,²⁷ R. Klingenberg,⁴⁶ T. Klingl,²³ T. Klioutchnikova,³² E.-E. Kluge,^{60a} P. Kluit,¹⁰⁹ S. Kluth,¹⁰³
 E. Kneringer,⁶⁵ E. B. F. G. Knoops,⁸⁸ A. Knue,¹⁰³ A. Kobayashi,¹⁵⁷ D. Kobayashi,¹⁵⁹ T. Kobayashi,¹⁵⁷ M. Kobel,⁴⁷
 M. Kocian,¹⁴⁵ P. Kodys,¹³¹ T. Koffas,³¹ E. Koffeman,¹⁰⁹ N. M. Köhler,¹⁰³ T. Koi,¹⁴⁵ M. Kolb,^{60b} I. Koletsou,⁵
 A. A. Komar,^{98,a} Y. Komori,¹⁵⁷ T. Kondo,⁶⁹ N. Kondrashova,^{36c} K. Köneke,⁵¹ A. C. König,¹⁰⁸ T. Kono,^{69,ee}
 R. Konoplich,^{112,ff} N. Konstantinidis,⁸¹ R. Kopeliantsky,⁶⁴ S. Koperny,^{41a} A. K. Kopp,⁵¹ K. Korcyl,⁴² K. Kordas,¹⁵⁶
 A. Korn,⁸¹ A. A. Korol,^{111,d} I. Korolkov,¹³ E. V. Korolkova,¹⁴¹ O. Kortner,¹⁰³ S. Kortner,¹⁰³ T. Kosek,¹³¹ V. V. Kostyukhin,²³
 A. Kotwal,⁴⁸ A. Koulouris,¹⁰ A. Kourkoumeli-Charalampidi,^{123a,123b} C. Kourkoumelis,⁹ E. Kourlitis,¹⁴¹ V. Kouskoura,²⁷
 A. B. Kowalewska,⁴² R. Kowalewski,¹⁷² T. Z. Kowalski,^{41a} C. Kozakai,¹⁵⁷ W. Kozanecki,¹³⁸ A. S. Kozhin,¹³²
 V. A. Kramarenko,¹⁰¹ G. Kramberger,⁷⁸ D. Krasnoperov,¹⁰⁰ M. W. Krasny,⁸³ A. Krasznahorkay,³² D. Krauss,¹⁰³
 J. A. Kremer,^{41a} J. Kretschmar,⁷⁷ K. Kreutzfeldt,⁵⁵ P. Krieger,¹⁶¹ K. Krizka,³³ K. Kroeninger,⁴⁶ H. Kroha,¹⁰³ J. Kroll,¹²⁹
 J. Kroll,¹²⁴ J. Kroseberg,²³ J. Krstic,¹⁴ U. Kruchonak,⁶⁸ H. Krüger,²³ N. Krumnack,⁶⁷ M. C. Kruse,⁴⁸ T. Kubota,⁹¹
 H. Kucuk,⁸¹ S. Kuday,^{4b} J. T. Kuechler,¹⁷⁸ S. Kuehn,³² A. Kugel,^{60a} F. Kuger,¹⁷⁷ T. Kuhl,⁴⁵ V. Kukhtin,⁶⁸ R. Kukla,⁸⁸
 Y. Kulchitsky,⁹⁵ S. Kuleshov,^{34b} Y. P. Kulinich,¹⁶⁹ M. Kuna,^{134a,134b} T. Kunigo,⁷¹ A. Kupco,¹²⁹ T. Kupfer,⁴⁶ O. Kuprash,¹⁵⁵
 H. Kurashige,⁷⁰ L. L. Kurchaninov,^{163a} Y. A. Kurochkin,⁹⁵ M. G. Kurth,^{35a} V. Kus,¹²⁹ E. S. Kuwertz,¹⁷² M. Kuze,¹⁵⁹
 J. Kvita,¹¹⁷ T. Kwan,¹⁷² D. Kyriazopoulos,¹⁴¹ A. La Rosa,¹⁰³ J. L. La Rosa Navarro,^{26d} L. La Rotonda,^{40a,40b}
 F. La Ruffa,^{40a,40b} C. Lacasta,¹⁷⁰ F. Lacava,^{134a,134b} J. Lacey,⁴⁵ H. Lacker,¹⁷ D. Lacour,⁸³ E. Ladygin,⁶⁸ R. Lafaye,⁵
 B. Laforge,⁸³ T. Lagouri,¹⁷⁹ S. Lai,⁵⁷ S. Lammers,⁶⁴ W. Lampl,⁷ E. Lançon,²⁷ U. Landgraf,⁵¹ M. P. J. Landon,⁷⁹
 M. C. Lanfermann,⁵² V. S. Lang,^{60a} J. C. Lange,¹³ R. J. Langenberg,³² A. J. Lankford,¹⁶⁶ F. Lanni,²⁷ K. Lantzsch,²³
 A. Lanza,^{123a} A. Lapertosa,^{53a,53b} S. Laplace,⁸³ J. F. Laporte,¹³⁸ T. Lari,^{94a} F. Lasagni Manghi,^{22a,22b} M. Lassnig,³²
 P. Laurelli,⁵⁰ W. Lavrijsen,¹⁶ A. T. Law,¹³⁹ P. Laycock,⁷⁷ T. Lazovich,⁵⁹ M. Lazzaroni,^{94a,94b} B. Le,⁹¹ O. Le Dortz,⁸³
 E. Le Guirriec,⁸⁸ E. P. Le Quilleuc,¹³⁸ M. LeBlanc,¹⁷² T. LeCompte,⁶ F. Ledroit-Guillon,⁵⁸ C. A. Lee,²⁷ G. R. Lee,^{133,gg}
 S. C. Lee,¹⁵³ L. Lee,⁵⁹ B. Lefebvre,⁹⁰ G. Lefebvre,⁸³ M. Lefebvre,¹⁷² F. Legger,¹⁰² C. Leggett,¹⁶ G. Lehmann Miotto,³²
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 B. Lenzi,³² R. Leone,⁷ S. Leone,^{126a,126b} C. Leonidopoulos,⁴⁹ G. Lerner,¹⁵¹ C. Leroy,⁹⁷ A. A. J. Lesage,¹³⁸ C. G. Lester,³⁰
 M. Levchenko,¹²⁵ J. Levêque,⁵ D. Levin,⁹² L. J. Levinson,¹⁷⁵ M. Levy,¹⁹ D. Lewis,⁷⁹ B. Li,^{36a,hh} Changqiao Li,^{36a} H. Li,¹⁵⁰
 L. Li,^{36c} Q. Li,^{35a} Q. Li,^{36a} S. Li,⁴⁸ X. Li,^{36c} Y. Li,¹⁴³ Z. Liang,^{35a} B. Liberti,^{135a} A. Liblong,¹⁶¹ K. Lie,^{62c} J. Liebal,²³
 W. Liebig,¹⁵ A. Limosani,¹⁵² S. C. Lin,¹⁸² T. H. Lin,⁸⁶ R. A. Linck,⁶⁴ B. E. Lindquist,¹⁵⁰ A. E. Lioni,⁵² E. Lipeles,¹²⁴
 A. Lipniacka,¹⁵ M. Lisovyi,^{60b} T. M. Liss,^{169,ii} A. Lister,¹⁷¹ A. M. Litke,¹³⁹ B. Liu,^{153,jj} H. Liu,⁹² H. Liu,²⁷ J. K. K. Liu,¹²²
 J. Liu,^{36b} J. B. Liu,^{36a} K. Liu,⁸⁸ L. Liu,¹⁶⁹ M. Liu,^{36a} Y. L. Liu,^{36a} Y. Liu,^{36a} M. Livan,^{123a,123b} A. Lleres,⁵⁸
 J. Llorente Merino,^{35a} S. L. Lloyd,⁷⁹ C. Y. Lo,^{62b} F. Lo Sterzo,¹⁵³ E. M. Lobodzinska,⁴⁵ P. Loch,⁷ F. K. Loebinger,⁸⁷
 A. Loesle,⁵¹ K. M. Loew,²⁵ A. Loginov,^{179,a} T. Lohse,¹⁷ K. Lohwasser,¹⁴¹ M. Lokajicek,¹²⁹ B. A. Long,²⁴ J. D. Long,¹⁶⁹
 R. E. Long,⁷⁵ L. Longo,^{76a,76b} K. A.Looper,¹¹³ J. A. Lopez,^{34b} D. Lopez Mateos,⁵⁹ I. Lopez Paz,¹³ A. Lopez Solis,⁸³
 J. Lorenz,¹⁰² N. Lorenzo Martinez,⁵ M. Losada,²¹ P. J. Lösel,¹⁰² X. Lou,^{35a} A. Lounis,¹¹⁹ J. Love,⁶ P. A. Love,⁷⁵ H. Lu,^{62a}
 N. Lu,⁹² Y. J. Lu,⁶³ H. J. Lubatti,¹⁴⁰ C. Luci,^{134a,134b} A. Lucotte,⁵⁸ C. Luedtke,⁵¹ F. Luehring,⁶⁴ W. Lukas,⁶⁵ L. Luminari,^{134a}
 O. Lundberg,^{148a,148b} B. Lund-Jensen,¹⁴⁹ M. S. Lutz,⁸⁹ P. M. Luzzi,⁸³ D. Lynn,²⁷ R. Lysak,¹²⁹ E. Lytken,⁸⁴ F. Lyu,^{35a}
 V. Lyubushkin,⁶⁸ H. Ma,²⁷ L. L. Ma,^{36b} Y. Ma,^{36b} G. Maccarrone,⁵⁰ A. Macchiolo,¹⁰³ C. M. Macdonald,¹⁴¹ B. Maček,⁷⁸
 J. Machado Miguens,^{124,128b} D. Madaffari,¹⁷⁰ R. Madar,³⁷ W. F. Mader,⁴⁷ A. Madsen,⁴⁵ J. Maeda,⁷⁰ S. Maeland,¹⁵
 T. Maeno,²⁷ A. S. Maevskiy,¹⁰¹ V. Magerl,⁵¹ J. Mahlstedt,¹⁰⁹ C. Maiani,¹¹⁹ C. Maidantchik,^{26a} A. A. Maier,¹⁰³ T. Maier,¹⁰²
 A. Maio,^{128a,128b,128d} O. Majersky,^{146a} S. Majewski,¹¹⁸ Y. Makida,⁶⁹ N. Makovec,¹¹⁹ B. Malaescu,⁸³ Pa. Malecki,⁴²
 V. P. Maleev,¹²⁵ F. Malek,⁵⁸ U. Mallik,⁶⁶ D. Malon,⁶ C. Malone,³⁰ S. Maltezos,¹⁰ S. Malyukov,³² J. Mamuzic,¹⁷⁰

G. Mancini,⁵⁰ I. Mandić,⁷⁸ J. Maneira,^{128a,128b} L. Manhaes de Andrade Filho,^{26b} J. Manjarres Ramos,⁴⁷ K. H. Mankinen,⁸⁴
A. Mann,¹⁰² A. Manousos,³² B. Mansoulie,¹³⁸ J. D. Mansour,^{35a} R. Mantifel,⁹⁰ M. Mantoani,⁵⁷ S. Manzoni,^{94a,94b}
L. Mapelli,³² G. Marceca,²⁹ L. March,⁵² L. Marchese,¹²² G. Marchiori,⁸³ M. Marcisovsky,¹²⁹ M. Marjanovic,³⁷
D. E. Marley,⁹² F. Marroquim,^{26a} S. P. Marsden,⁸⁷ Z. Marshall,¹⁶ M. U. F. Martensson,¹⁶⁸ S. Marti-Garcia,¹⁷⁰ C. B. Martin,¹¹³
T. A. Martin,¹⁷³ V. J. Martin,⁴⁹ B. Martin dit Latour,¹⁵ M. Martinez,^{13,w} V. I. Martinez Outschoorn,¹⁶⁹ S. Martin-Haugh,¹³³
V. S. Martoiu,^{28b} A. C. Martyniuk,⁸¹ A. Marzin,³² L. Masetti,⁸⁶ T. Mashimo,¹⁵⁷ R. Mashinistov,⁹⁸ J. Masik,⁸⁷
A. L. Maslennikov,^{111,d} L. Massa,^{135a,135b} P. Mastrandrea,⁵ A. Mastroberardino,^{40a,40b} T. Masubuchi,¹⁵⁷ P. Mättig,¹⁷⁸
J. Maurer,^{28b} S. J. Maxfield,⁷⁷ D. A. Maximov,^{111,d} R. Mazini,¹⁵³ I. Maznas,¹⁵⁶ S. M. Mazza,^{94a,94b} N. C. Mc Fadden,¹⁰⁷
G. Mc Goldrick,¹⁶¹ S. P. Mc Kee,⁹² A. McCarn,⁹² R. L. McCarthy,¹⁵⁰ T. G. McCarthy,¹⁰³ L. I. McClymont,⁸¹
E. F. McDonald,⁹¹ J. A. Mcfayden,⁸¹ G. Mchedlidze,⁵⁷ S. J. McMahon,¹³³ P. C. McNamara,⁹¹ R. A. McPherson,^{172,p}
S. Meehan,¹⁴⁰ T. J. Megy,⁵¹ S. Mehlhase,¹⁰² A. Mehta,⁷⁷ T. Meideck,⁵⁸ K. Meier,^{60a} B. Meirose,⁴⁴ D. Melini,^{170,kk}
B. R. Mellado Garcia,^{147c} J. D. Mellenthin,⁵⁷ M. Melo,^{146a} F. Meloni,¹⁸ A. Melzer,²³ S. B. Menary,⁸⁷ L. Meng,⁷⁷
X. T. Meng,⁹² A. Mengarelli,^{22a,22b} S. Menke,¹⁰³ E. Meoni,^{40a,40b} S. Mergelmeyer,¹⁷ P. Mermod,⁵² L. Merola,^{106a,106b}
C. Meroni,^{94a} F. S. Merritt,³³ A. Messina,^{134a,134b} J. Metcalfe,⁶ A. S. Mete,¹⁶⁶ C. Meyer,¹²⁴ J-P. Meyer,¹³⁸ J. Meyer,¹⁰⁹
H. Meyer Zu Theenhausen,^{60a} F. Miano,¹⁵¹ R. P. Middleton,¹³³ S. Miglioranzi,^{53a,53b} L. Mijović,⁴⁹ G. Mikenberg,¹⁷⁵
M. Mikestikova,¹²⁹ M. Mikuž,⁷⁸ M. Milesi,⁹¹ A. Milic,¹⁶¹ D. W. Miller,³³ C. Mills,⁴⁹ A. Milov,¹⁷⁵ D. A. Milstead,^{148a,148b}
A. A. Minaenko,¹³² Y. Minami,¹⁵⁷ I. A. Minashvili,⁶⁸ A. I. Mincer,¹¹² B. Mindur,^{41a} M. Mineev,⁶⁸ Y. Minegishi,¹⁵⁷
Y. Ming,¹⁷⁶ L. M. Mir,¹³ K. P. Mistry,¹²⁴ T. Mitani,¹⁷⁴ J. Mitrevski,¹⁰² V. A. Mitsou,¹⁷⁰ A. Miucci,¹⁸ P. S. Miyagawa,¹⁴¹
A. Mizukami,⁶⁹ J. U. Mjörnmark,⁸⁴ T. Mkrtchyan,¹⁸⁰ M. Mlynarikova,¹³¹ T. Moa,^{148a,148b} K. Mochizuki,⁹⁷ P. Mogg,⁵¹
S. Mohapatra,³⁸ S. Molander,^{148a,148b} R. Moles-Valls,²³ R. Monden,⁷¹ M. C. Mondragon,⁹³ K. Mönig,⁴⁵ J. Monk,³⁹
E. Monnier,⁸⁸ A. Montalbano,¹⁵⁰ J. Montejo Berlingen,³² F. Monticelli,⁷⁴ S. Monzani,^{94a,94b} R. W. Moore,³ N. Morange,¹¹⁹
D. Moreno,²¹ M. Moreno Llácer,³² P. Morettini,^{53a} S. Morgenstern,³² D. Mori,¹⁴⁴ T. Mori,¹⁵⁷ M. Morii,⁵⁹ M. Morinaga,¹⁵⁷
V. Morisbak,¹²¹ A. K. Morley,³² G. Mornacchi,³² J. D. Morris,⁷⁹ L. Morvaj,¹⁵⁰ P. Moschovakos,¹⁰ M. Mosidze,^{54b}
H. J. Moss,¹⁴¹ J. Moss,^{145,ll} K. Motohashi,¹⁵⁹ R. Mount,¹⁴⁵ E. Mountricha,²⁷ E. J. W. Moyse,⁸⁹ S. Muanza,⁸⁸ F. Mueller,¹⁰³
J. Mueller,¹²⁷ R. S. P. Mueller,¹⁰² D. Muenstermann,⁷⁵ P. Mullen,⁵⁶ G. A. Mullier,¹⁸ F. J. Munoz Sanchez,⁸⁷
W. J. Murray,^{173,133} H. Musheghyan,³² M. Muškinja,⁷⁸ A. G. Myagkov,^{132,mm} M. Myska,¹³⁰ B. P. Nachman,¹⁶
O. Nackenhorst,⁵² K. Nagai,¹²² R. Nagai,^{69,ee} K. Nagano,⁶⁹ Y. Nagasaka,⁶¹ K. Nagata,¹⁶⁴ M. Nagel,⁵¹ E. Nagy,⁸⁸
A. M. Nairz,³² Y. Nakahama,¹⁰⁵ K. Nakamura,⁶⁹ T. Nakamura,¹⁵⁷ I. Nakano,¹¹⁴ R. F. Naranjo Garcia,⁴⁵ R. Narayan,¹¹
D. I. Narrias Villar,^{60a} I. Naryshkin,¹²⁵ T. Naumann,⁴⁵ G. Navarro,²¹ R. Nayyar,⁷ H. A. Neal,⁹² P. Yu. Nechaeva,⁹⁸
T. J. Neep,¹³⁸ A. Negri,^{123a,123b} M. Negrini,^{22a} S. Nektarijevic,¹⁰⁸ C. Nellist,¹¹⁹ A. Nelson,¹⁶⁶ M. E. Nelson,¹²²
S. Nemecek,¹²⁹ P. Nemethy,¹¹² M. Nessi,^{32,nn} M. S. Neubauer,¹⁶⁹ M. Neumann,¹⁷⁸ P. R. Newman,¹⁹ T. Y. Ng,^{62c}
T. Nguyen Manh,⁹⁷ R. B. Nickerson,¹²² R. Nicolaidou,¹³⁸ J. Nielsen,¹³⁹ V. Nikolaenko,^{132,mm} I. Nikolic-Audit,⁸³
K. Nikolopoulos,¹⁹ J. K. Nilsen,¹²¹ P. Nilsson,²⁷ Y. Ninomiya,¹⁵⁷ A. Nisati,^{134a} N. Nishu,^{35c} R. Nisius,¹⁰³ I. Nitsche,⁴⁶
T. Nitta,¹⁷⁴ T. Nobe,¹⁵⁷ Y. Noguchi,⁷¹ M. Nomachi,¹²⁰ I. Nomidis,³¹ M. A. Nomura,²⁷ T. Nooney,⁷⁹ M. Nordberg,³²
N. Norjoharuddeen,¹²² O. Novgorodova,⁴⁷ M. Nozaki,⁶⁹ L. Nozka,¹¹⁷ K. Ntekas,¹⁶⁶ E. Nurse,⁸¹ F. Nuti,⁹¹ K. O'connor,²⁵
D. C. O'Neil,¹⁴⁴ A. A. O'Rourke,⁴⁵ V. O'Shea,⁵⁶ F. G. Oakham,^{31,e} H. Oberlack,¹⁰³ T. Obermann,²³ J. Ocariz,⁸³ A. Ochi,⁷⁰
I. Ochoa,³⁸ J. P. Ochoa-Ricoux,^{34a} S. Oda,⁷³ S. Odaka,⁶⁹ A. Oh,⁸⁷ S. H. Oh,⁴⁸ C. C. Ohm,¹⁶ H. Ohman,¹⁶⁸ H. Oide,^{53a,53b}
H. Okawa,¹⁶⁴ Y. Okumura,¹⁵⁷ T. Okuyama,⁶⁹ A. Olariu,^{28b} L. F. Oleiro Seabra,^{128a} S. A. Olivares Pino,^{34a}
D. Oliveira Damazio,²⁷ A. Olszewski,⁴² J. Olszowska,⁴² A. Onofre,^{128a,128e} K. Onogi,¹⁰⁵ P. U. E. Onyisi,^{11,aa} H. Oppen,¹²¹
M. J. Oreglia,³³ Y. Oren,¹⁵⁵ D. Orestano,^{136a,136b} N. Orlando,^{62b} R. S. Orr,¹⁶¹ B. Osculati,^{53a,53b,a} R. Ospanov,^{36a}
G. Otero y Garzon,²⁹ H. Otono,⁷³ M. Ouchrif,^{137d} F. Ould-Saada,¹²¹ A. Ouraou,¹³⁸ K. P. Oussoren,¹⁰⁹ Q. Ouyang,^{35a}
M. Owen,⁵⁶ R. E. Owen,¹⁹ V. E. Ozcan,^{20a} N. Ozturk,⁸ K. Pachal,¹⁴⁴ A. Pacheco Pages,¹³ L. Pacheco Rodriguez,¹³⁸
C. Padilla Aranda,¹³ S. Pagan Griso,¹⁶ M. Paganini,¹⁷⁹ F. Paige,²⁷ G. Palacino,⁶⁴ S. Palazzo,^{40a,40b} S. Palestini,³² M. Palka,^{41b}
D. Pallin,³⁷ E. St. Panagiotopoulou,¹⁰ I. Panagoulas,¹⁰ C. E. Pandini,^{126a,126b} J. G. Panduro Vazquez,⁸⁰ P. Pani,³²
S. Panitkin,²⁷ D. Pantea,^{28b} L. Paolozzi,⁵² Th. D. Papadopoulou,¹⁰ K. Papageorgiou,^{9,t} A. Paramonov,⁶
D. Paredes Hernandez,¹⁷⁹ A. J. Parker,⁷⁵ M. A. Parker,³⁰ K. A. Parker,⁴⁵ F. Parodi,^{53a,53b} J. A. Parsons,³⁸ U. Parzefall,⁵¹
V. R. Pascuzzi,¹⁶¹ J. M. Pasner,¹³⁹ E. Pasqualucci,^{134a} S. Passaggio,^{53a} Fr. Pastore,⁸⁰ S. Pataraja,⁸⁶ J. R. Pater,⁸⁷ T. Pauly,³²
B. Pearson,¹⁰³ S. Pedraza Lopez,¹⁷⁰ R. Pedro,^{128a,128b} S. V. Peleganchuk,^{111,d} O. Penc,¹²⁹ C. Peng,^{35a} H. Peng,^{36a} J. Penwell,⁶⁴
B. S. Peralva,^{26b} M. M. Perego,¹³⁸ D. V. Perepelitsa,²⁷ F. Peri,¹⁷ L. Perini,^{94a,94b} H. Pernegger,³² S. Perrella,^{106a,106b}

R. Peschke,⁴⁵ V. D. Peshekhonov,^{68,a} K. Peters,⁴⁵ R. F. Y. Peters,⁸⁷ B. A. Petersen,³² T. C. Petersen,³⁹ E. Petit,⁵⁸ A. Petridis,¹ C. Petridou,¹⁵⁶ P. Petroff,¹¹⁹ E. Petrolo,^{134a} M. Petrov,¹²² F. Petrucci,^{136a,136b} N. E. Pettersson,⁸⁹ A. Peyaud,¹³⁸ R. Pezoa,^{34b} F. H. Phillips,⁹³ P. W. Phillips,¹³³ G. Piacquadio,¹⁵⁰ E. Pianori,¹⁷³ A. Picazio,⁸⁹ E. Piccaro,⁷⁹ M. A. Pickering,¹²² R. Piegaia,²⁹ J. E. Pilcher,³³ A. D. Pilkington,⁸⁷ A. W. J. Pin,⁸⁷ M. Pinamonti,^{135a,135b} J. L. Pinfeld,³ H. Pirumov,⁴⁵ M. Pitt,¹⁷⁵ L. Plazak,^{146a} M.-A. Pleier,²⁷ V. Pleskot,⁸⁶ E. Plotnikova,⁶⁸ D. Pluth,⁶⁷ P. Podberezko,¹¹¹ R. Poettgen,⁸⁴ R. Poggi,^{123a,123b} L. Poggioli,¹¹⁹ I. Pogrebnyak,⁹³ D. Pohl,²³ G. Polesello,^{123a} A. Poley,⁴⁵ A. Policicchio,^{40a,40b} R. Polifka,³² A. Polini,^{22a} C. S. Pollard,⁵⁶ V. Polychronakos,²⁷ K. Pommès,³² D. Ponomarenko,¹⁰⁰ L. Pontecorvo,^{134a} G. A. Popeneciu,^{28d} S. Pospisil,¹³⁰ K. Potamianos,¹⁶ I. N. Potrap,⁶⁸ C. J. Potter,³⁰ T. Poulsen,⁸⁴ J. Poveda,³² M. E. Pozo Astigarraga,³² P. Pralavorio,⁸⁸ A. Pranko,¹⁶ S. Prell,⁶⁷ D. Price,⁸⁷ M. Primavera,^{76a} S. Prince,⁹⁰ N. Proklova,¹⁰⁰ K. Prokofiev,^{62c} F. Prokoshin,^{34b} S. Protopopescu,²⁷ J. Proudfoot,⁶ M. Przybycien,^{41a} A. Puri,¹⁶⁹ P. Puzo,¹¹⁹ J. Qian,⁹² G. Qin,⁵⁶ Y. Qin,⁸⁷ A. Quadt,⁵⁷ M. Queitsch-Maitland,⁴⁵ D. Quilty,⁵⁶ S. Raddum,¹²¹ V. Radeka,²⁷ V. Radescu,¹²² S. K. Radhakrishnan,¹⁵⁰ P. Radloff,¹¹⁸ P. Rados,⁹¹ F. Ragusa,^{94a,94b} G. Rahal,¹⁸¹ J. A. Raine,⁸⁷ S. Rajagopalan,²⁷ C. Rangel-Smith,¹⁶⁸ T. Rashid,¹¹⁹ S. Raspopov,⁵ M. G. Ratti,^{94a,94b} D. M. Rauch,⁴⁵ F. Rauscher,¹⁰² S. Rave,⁸⁶ I. Ravinovich,¹⁷⁵ J. H. Rawling,⁸⁷ M. Raymond,³² A. L. Read,¹²¹ N. P. Readioff,⁵⁸ M. Reale,^{76a,76b} D. M. Rebuffi,^{123a,123b} A. Redelbach,¹⁷⁷ G. Redlinger,²⁷ R. Reece,¹³⁹ R. G. Reed,^{147c} K. Reeves,⁴⁴ L. Rehnisch,¹⁷ J. Reichert,¹²⁴ A. Reiss,⁸⁶ C. Rembser,³² H. Ren,^{35a} M. Rescigno,^{134a} S. Resconi,^{94a} E. D. Resseguie,¹²⁴ S. Rettie,¹⁷¹ E. Reynolds,¹⁹ O. L. Rezanova,^{111,d} P. Reznicek,¹³¹ R. Rezvani,⁹⁷ R. Richter,¹⁰³ S. Richter,⁸¹ E. Richter-Was,^{41b} O. Ricken,²³ M. Ridel,⁸³ P. Rieck,¹⁰³ C. J. Riegel,¹⁷⁸ J. Rieger,⁵⁷ O. Rifki,¹¹⁵ M. Rijssenbeek,¹⁵⁰ A. Rimoldi,^{123a,123b} M. Rimoldi,¹⁸ L. Rinaldi,^{22a} G. Ripellino,¹⁴⁹ B. Ristić,³² E. Ritsch,³² I. Riu,¹³ F. Rizatdinova,¹¹⁶ E. Rizvi,⁷⁹ C. Rizzi,¹³ R. T. Roberts,⁸⁷ S. H. Robertson,^{90,p} A. Robichaud-Veronneau,⁹⁰ D. Robinson,³⁰ J. E. M. Robinson,⁴⁵ A. Robson,⁵⁶ E. Rocco,⁸⁶ C. Roda,^{126a,126b} Y. Rodina,^{88,oo} S. Rodriguez Bosca,¹⁷⁰ A. Rodriguez Perez,¹³ D. Rodriguez Rodriguez,¹⁷⁰ S. Roe,³² C. S. Rogan,⁵⁹ O. Røhne,¹²¹ J. Roloff,⁵⁹ A. Romaniouk,¹⁰⁰ M. Romano,^{22a,22b} S. M. Romano Saez,³⁷ E. Romero Adam,¹⁷⁰ N. Rompotis,⁷⁷ M. Ronzani,⁵¹ L. Roos,⁸³ S. Rosati,^{134a} K. Rosbach,⁵¹ P. Rose,¹³⁹ N.-A. Rosien,⁵⁷ E. Rossi,^{106a,106b} L. P. Rossi,^{53a} J. H. N. Rosten,³⁰ R. Rosten,¹⁴⁰ M. Rotaru,^{28b} J. Rothberg,¹⁴⁰ D. Rousseau,¹¹⁹ A. Rozanov,⁸⁸ Y. Rozen,¹⁵⁴ X. Ruan,^{147c} F. Rubbo,¹⁴⁵ F. Rühr,⁵¹ A. Ruiz-Martinez,³¹ Z. Rurikova,⁵¹ N. A. Rusakovich,⁶⁸ H. L. Russell,⁹⁰ J. P. Rutherford,⁷ N. Ruthmann,³² Y. F. Ryabov,¹²⁵ M. Rybar,¹⁶⁹ G. Rybkin,¹¹⁹ S. Ryu,⁶ A. Ryzhov,¹³² G. F. Rzehorz,⁵⁷ A. F. Saavedra,¹⁵² G. Sabato,¹⁰⁹ S. Sacerdoti,²⁹ H. F.-W. Sadrozinski,¹³⁹ R. Sadykov,⁶⁸ F. Safai Tehrani,^{134a} P. Saha,¹¹⁰ M. Sahinsoy,^{60a} M. Saimpert,⁴⁵ M. Saito,¹⁵⁷ T. Saito,¹⁵⁷ H. Sakamoto,¹⁵⁷ Y. Sakurai,¹⁷⁴ G. Salamanna,^{136a,136b} J. E. Salazar Loyola,^{34b} D. Salek,¹⁰⁹ P. H. Sales De Bruin,¹⁶⁸ D. Salihagic,¹⁰³ A. Salnikov,¹⁴⁵ J. Salt,¹⁷⁰ D. Salvatore,^{40a,40b} F. Salvatore,¹⁵¹ A. Salvucci,^{62a,62b,62c} A. Salzburger,³² D. Sammel,⁵¹ D. Sampsonidis,¹⁵⁶ D. Sampsonidou,¹⁵⁶ J. Sánchez,¹⁷⁰ V. Sanchez Martinez,¹⁷⁰ A. Sanchez Pineda,^{167a,167c} H. Sandaker,¹²¹ R. L. Sandbach,⁷⁹ C. O. Sander,⁴⁵ M. Sandhoff,¹⁷⁸ C. Sandoval,²¹ D. P. C. Sankey,¹³³ M. Sannino,^{53a,53b} Y. Sano,¹⁰⁵ A. Sansoni,⁵⁰ C. Santoni,³⁷ H. Santos,^{128a} I. Santoyo Castillo,¹⁵¹ A. Sapronov,⁶⁸ J. G. Saraiva,^{128a,128d} B. Sarrazin,²³ O. Sasaki,⁶⁹ K. Sato,¹⁶⁴ E. Sauvan,⁵ G. Savage,⁸⁰ P. Savard,^{161,e} N. Savić,¹⁰³ C. Sawyer,¹³³ L. Sawyer,^{82,v} J. Saxon,³³ C. Sbarra,^{22a} A. Sbrizzi,^{22a,22b} T. Scanlon,⁸¹ D. A. Scannicchio,¹⁶⁶ M. Scarcella,¹⁵² J. Schaarschmidt,¹⁴⁰ P. Schacht,¹⁰³ B. M. Schachtner,¹⁰² D. Schaefer,³² L. Schaefer,¹²⁴ R. Schaefer,⁴⁵ J. Schaeffer,⁸⁶ S. Schaepe,²³ S. Schaezel,^{60b} U. Schäfer,⁸⁶ A. C. Schaffer,¹¹⁹ D. Schaile,¹⁰² R. D. Schamberger,¹⁵⁰ V. A. Schegelsky,¹²⁵ D. Scheirich,¹³¹ M. Schernau,¹⁶⁶ C. Schiavi,^{53a,53b} S. Schier,¹³⁹ L. K. Schildgen,²³ C. Schillo,⁵¹ M. Schioppa,^{40a,40b} S. Schlenker,³² K. R. Schmidt-Sommerfeld,¹⁰³ K. Schmieden,³² C. Schmitt,⁸⁶ S. Schmitt,⁴⁵ S. Schmitz,⁸⁶ U. Schnoor,⁵¹ L. Schoeffel,¹³⁸ A. Schoening,^{60b} B. D. Schoenrock,⁹³ E. Schopf,²³ M. Schott,⁸⁶ J. F. P. Schouwenberg,¹⁰⁸ J. Schovancova,³² S. Schramm,⁵² N. Schuh,⁸⁶ A. Schulte,⁸⁶ M. J. Schultens,²³ H.-C. Schultz-Coulon,^{60a} H. Schulz,¹⁷ M. Schumacher,⁵¹ B. A. Schumm,¹³⁹ Ph. Schune,¹³⁸ A. Schwartzman,¹⁴⁵ T. A. Schwarz,⁹² H. Schweiger,⁸⁷ Ph. Schwemling,¹³⁸ R. Schwienhorst,⁹³ J. Schwindling,¹³⁸ A. Sciandra,²³ G. Sciolla,²⁵ M. Scornajenghi,^{40a,40b} F. Scuri,^{126a,126b} F. Scutti,⁹¹ J. Searcy,⁹² P. Seema,²³ S. C. Seidel,¹⁰⁷ A. Seiden,¹³⁹ J. M. Seixas,^{26a} G. Sekhniaidze,^{106a} K. Sekhon,⁹² S. J. Sekula,⁴³ N. Semprini-Cesari,^{22a,22b} S. Senkin,³⁷ C. Serfon,¹²¹ L. Serin,¹¹⁹ L. Serkin,^{167a,167b} M. Sessa,^{136a,136b} R. Seuster,¹⁷² H. Severini,¹¹⁵ T. Sfiligoj,⁷⁸ F. Sforza,³² A. Sfyrla,⁵² E. Shabalina,⁵⁷ N. W. Shaikh,^{148a,148b} L. Y. Shan,^{35a} R. Shang,¹⁶⁹ J. T. Shank,²⁴ M. Shapiro,¹⁶ P. B. Shatalov,⁹⁹ K. Shaw,^{167a,167b} S. M. Shaw,⁸⁷ A. Shcherbakova,^{148a,148b} C. Y. Shehu,¹⁵¹ Y. Shen,¹¹⁵ N. Sherafati,³¹ P. Sherwood,⁸¹ L. Shi,^{153,pp} S. Shimizu,⁷⁰ C. O. Shimmmin,¹⁷⁹ M. Shimojima,¹⁰⁴ I. P. J. Shipsey,¹²² S. Shirabe,⁷³ M. Shiyakova,^{68,qq} J. Shlomi,¹⁷⁵ A. Shmeleva,⁹⁸ D. Shoaleh Saadi,⁹⁷ M. J. Shochet,³³ S. Shojaii,^{94a} D. R. Shope,¹¹⁵ S. Shrestha,¹¹³ E. Shulga,¹⁰⁰ M. A. Shupe,⁷ P. Sicho,¹²⁹

A. M. Sickles,¹⁶⁹ P. E. Sidebo,¹⁴⁹ E. Sideras Haddad,^{147c} O. Sidiropoulou,¹⁷⁷ A. Sidoti,^{22a,22b} F. Siegert,⁴⁷ Dj. Sijacki,¹⁴ J. Silva,^{128a,128d} S. B. Silverstein,^{148a} V. Simak,¹³⁰ Lj. Simic,¹⁴ S. Simion,¹¹⁹ E. Simioni,⁸⁶ B. Simmons,⁸¹ M. Simon,⁸⁶ P. Sinervo,¹⁶¹ N. B. Sinev,¹¹⁸ M. Sioli,^{22a,22b} G. Siragusa,¹⁷⁷ I. Siral,⁹² S. Yu. Sivoklov,¹⁰¹ J. Sjölin,^{148a,148b} M. B. Skinner,⁷⁵ P. Skubic,¹¹⁵ M. Slater,¹⁹ T. Slavicek,¹³⁰ M. Slawinska,⁴² K. Sliwa,¹⁶⁵ R. Slovak,¹³¹ V. Smakhtin,¹⁷⁵ B. H. Smart,⁵ J. Smiesko,^{146a} N. Smirnov,¹⁰⁰ S. Yu. Smirnov,¹⁰⁰ Y. Smirnov,¹⁰⁰ L. N. Smirnova,^{101,rr} O. Smirnova,⁸⁴ J. W. Smith,⁵⁷ M. N. K. Smith,³⁸ R. W. Smith,³⁸ M. Smizanska,⁷⁵ K. Smolek,¹³⁰ A. A. Snesarev,⁹⁸ I. M. Snyder,¹¹⁸ S. Snyder,²⁷ R. Sobie,^{172,p} F. Socher,⁴⁷ A. Soffer,¹⁵⁵ A. Søggaard,⁴⁹ D. A. Soh,¹⁵³ G. Sokhrannyi,⁷⁸ C. A. Solans Sanchez,³² M. Solar,¹³⁰ E. Yu. Soldatov,¹⁰⁰ U. Soldevila,¹⁷⁰ A. A. Solodkov,¹³² A. Soloshenko,⁶⁸ O. V. Solovyanov,¹³² V. Solovyev,¹²⁵ P. Sommer,⁵¹ H. Son,¹⁶⁵ A. Sopczak,¹³⁰ D. Sosa,^{60b} C. L. Sotiropoulou,^{126a,126b} R. Soualah,^{167a,167c} A. M. Soukharev,^{111,d} D. South,⁴⁵ B. C. Sowden,⁸⁰ S. Spagnolo,^{76a,76b} M. Spalla,^{126a,126b} M. Spangenberg,¹⁷³ F. Spanò,⁸⁰ D. Sperlich,¹⁷ F. Spettel,¹⁰³ T. M. Spieker,^{60a} R. Spighi,^{22a} G. Spigo,³² L. A. Spiller,⁹¹ M. Spousta,¹³¹ R. D. St. Denis,^{56a} A. Stabile,^{94a} R. Stamen,^{60a} S. Stamm,¹⁷ E. Stanecka,⁴² R. W. Stanek,⁶ C. Stanescu,^{136a} M. M. Stanitzki,⁴⁵ B. S. Stapf,¹⁰⁹ S. Stapnes,¹²¹ E. A. Starchenko,¹³² G. H. Stark,³³ J. Stark,⁵⁸ S. H. Stark,³⁹ P. Staroba,¹²⁹ P. Starovoitov,^{60a} S. Stärz,³² R. Staszewski,⁴² P. Steinberg,²⁷ B. Stelzer,¹⁴⁴ H. J. Stelzer,³² O. Stelzer-Chilton,^{163a} H. Stenzel,⁵⁵ G. A. Stewart,⁵⁶ M. C. Stockton,¹¹⁸ M. Stoebe,⁹⁰ G. Stoica,^{28b} P. Stolte,⁵⁷ S. Stonjek,¹⁰³ A. R. Stradling,⁸ A. Straessner,⁴⁷ M. E. Stramaglia,¹⁸ J. Strandberg,¹⁴⁹ S. Strandberg,^{148a,148b} M. Strauss,¹¹⁵ P. Strizenec,^{146b} R. Ströhmer,¹⁷⁷ D. M. Strom,¹¹⁸ R. Stroynowski,⁴³ A. Strubig,⁴⁹ S. A. Stucci,²⁷ B. Stugu,¹⁵ N. A. Styles,⁴⁵ D. Su,¹⁴⁵ J. Su,¹²⁷ S. Suchek,^{60a} Y. Sugaya,¹²⁰ M. Suk,¹³⁰ V. V. Sulin,⁹⁸ DMS Sultan,^{162a,162b} S. Sultansoy,^{4c} T. Sumida,⁷¹ S. Sun,⁵⁹ X. Sun,³ K. Suruliz,¹⁵¹ C. J. E. Suster,¹⁵² M. R. Sutton,¹⁵¹ S. Suzuki,⁶⁹ M. Svatos,¹²⁹ M. Swiatlowski,³³ S. P. Swift,² I. Sykora,^{146a} T. Sykora,¹³¹ D. Ta,⁵¹ K. Tackmann,⁴⁵ J. Taenzer,¹⁵⁵ A. Taffard,¹⁶⁶ R. Tafirout,^{163a} E. Tahirovic,⁷⁹ N. Taiblum,¹⁵⁵ H. Takai,²⁷ R. Takashima,⁷² E. H. Takasugi,¹⁰³ T. Takeshita,¹⁴² Y. Takubo,⁶⁹ M. Talby,⁸⁸ A. A. Talyshev,^{111,d} J. Tanaka,¹⁵⁷ M. Tanaka,¹⁵⁹ R. Tanaka,¹¹⁹ S. Tanaka,⁶⁹ R. Tanioka,⁷⁰ B. B. Tannenwald,¹¹³ S. Tapia Araya,^{34b} S. Tapprogge,⁸⁶ S. Tarem,¹⁵⁴ G. F. Tartarelli,^{94a} P. Tas,¹³¹ M. Tasevsky,¹²⁹ T. Tashiro,⁷¹ E. Tassi,^{40a,40b} A. Tavares Delgado,^{128a,128b} Y. Tayalati,^{137e} A. C. Taylor,¹⁰⁷ G. N. Taylor,⁹¹ P. T. E. Taylor,⁹¹ W. Taylor,^{163b} P. Teixeira-Dias,⁸⁰ D. Temple,¹⁴⁴ H. Ten Kate,³² P. K. Teng,¹⁵³ J. J. Teoh,¹²⁰ F. Tepel,¹⁷⁸ S. Terada,⁶⁹ K. Terashi,¹⁵⁷ J. Terron,⁸⁵ S. Terzo,¹³ M. Testa,⁵⁰ R. J. Teuscher,^{161,p} T. Thevenaux-Pelzer,⁸⁸ F. Thiele,³⁹ J. P. Thomas,¹⁹ J. Thomas-Wilsker,⁸⁰ P. D. Thompson,¹⁹ A. S. Thompson,⁵⁶ L. A. Thomsen,¹⁷⁹ E. Thomson,¹²⁴ M. J. Tibbetts,¹⁶ R. E. Tisce Torres,⁸⁸ V. O. Tikhomirov,^{98,ss} Yu. A. Tikhonov,^{111,d} S. Timoshenko,¹⁰⁰ P. Tipton,¹⁷⁹ S. Tisserant,⁸⁸ K. Todome,¹⁵⁹ S. Todorova-Nova,⁵ S. Todt,⁴⁷ J. Tojo,⁷³ S. Tokár,^{146a} K. Tokushuku,⁶⁹ E. Tolley,⁵⁹ L. Tomlinson,⁸⁷ M. Tomoto,¹⁰⁵ L. Tompkins,^{145,tt} K. Toms,¹⁰⁷ B. Tong,⁵⁹ P. Tornambe,⁵¹ E. Torrence,¹¹⁸ H. Torres,¹⁴⁴ E. Torró Pastor,¹⁴⁰ J. Toth,^{88,uu} F. Touchard,⁸⁸ D. R. Tovey,¹⁴¹ C. J. Treado,¹¹² T. Trefzger,¹⁷⁷ F. Tresoldi,¹⁵¹ A. Tricoli,²⁷ I. M. Trigger,^{163a} S. Trincaz-Duvoid,⁸³ M. F. Tripiana,¹³ W. Trischuk,¹⁶¹ B. Trocmé,⁵⁸ A. Trofymov,⁴⁵ C. Troncon,^{94a} M. Trotter-McDonald,¹⁶ M. Trovatelli,¹⁷² L. Truong,^{147b} M. Trzebinski,⁴² A. Trzupek,⁴² K. W. Tsang,^{62a} J. C-L. Tseng,¹²² P. V. Tsiarehka,⁹⁵ G. Tsipolitis,¹⁰ N. Tsirintanis,⁹ S. Tsiskaridze,¹³ V. Tsiskaridze,⁵¹ E. G. Tskhadadze,^{54a} K. M. Tsui,^{62a} I. I. Tsukerman,⁹⁹ V. Tsulaia,¹⁶ S. Tsuno,⁶⁹ D. Tsybychev,¹⁵⁰ Y. Tu,^{62b} A. Tudorache,^{28b} V. Tudorache,^{28b} T. T. Tulbure,^{28a} A. N. Tuna,⁵⁹ S. A. Tuppuri,^{22a,22b} S. Turchikhin,⁶⁸ D. Turgeman,¹⁷⁵ I. Turk Cakir,^{4b,vv} R. Turra,^{94a} P. M. Tuts,³⁸ G. Uccielli,^{22a,22b} I. Ueda,⁶⁹ M. Ughetto,^{148a,148b} F. Ukegawa,¹⁶⁴ G. Unal,³² A. Undrus,²⁷ G. Unel,¹⁶⁶ F. C. Ungaro,⁹¹ Y. Unno,⁶⁹ C. Unverdorben,¹⁰² J. Urban,^{146b} P. Urquijo,⁹¹ P. Urrejola,⁸⁶ G. Usai,⁸ J. Usui,⁶⁹ L. Vacavant,⁸⁸ V. Vacek,¹³⁰ B. Vachon,⁹⁰ K. O. H. Vadla,¹²¹ A. Vaidya,⁸¹ C. Valderanis,¹⁰² E. Valdes Santurio,^{148a,148b} M. Valente,⁵² S. Valentinetti,^{22a,22b} A. Valero,¹⁷⁰ L. Valéry,¹³ S. Valkar,¹³¹ A. Vallier,⁵ J. A. Valls Ferrer,¹⁷⁰ W. Van Den Wollenberg,¹⁰⁹ H. van der Graaf,¹⁰⁹ P. van Gemmeren,⁶ J. Van Nieuwkoop,¹⁴⁴ I. van Vulpen,¹⁰⁹ M. C. van Woerden,¹⁰⁹ M. Vanadia,^{135a,135b} W. Vandelli,³² A. Vaniachine,¹⁶⁰ P. Vankov,¹⁰⁹ G. Vardanyan,¹⁸⁰ R. Vari,^{134a} E. W. Varnes,⁷ C. Varni,^{53a,53b} T. Varol,⁴³ D. Varouchas,¹¹⁹ A. Vartapetian,⁸ K. E. Varvell,¹⁵² J. G. Vasquez,¹⁷⁹ G. A. Vasquez,^{34b} F. Vazeille,³⁷ T. Vazquez Schroeder,⁹⁰ J. Veatch,⁵⁷ V. Veeraraghavan,⁷ L. M. Veloce,¹⁶¹ F. Veloso,^{128a,128c} S. Veneziano,^{134a} A. Ventura,^{76a,76b} M. Venturi,¹⁷² N. Venturi,³² A. Venturini,²⁵ V. Vercesi,^{123a} M. Verducci,^{136a,136b} W. Verkerke,¹⁰⁹ A. T. Vermeulen,¹⁰⁹ J. C. Vermeulen,¹⁰⁹ M. C. Vetterli,^{144,e} N. Viaux Maira,^{34b} O. Viazlo,⁸⁴ I. Vichou,^{169,a} T. Vickey,¹⁴¹ O. E. Vickey Boeriu,¹⁴¹ G. H. A. Viehhauser,¹²² S. Viel,¹⁶ L. Vigani,¹²² M. Villa,^{22a,22b} M. Villaplana Perez,^{94a,94b} E. Vilucchi,⁵⁰ M. G. Vincker,³¹ V. B. Vinogradov,⁶⁸ A. Vishwakarma,⁴⁵ C. Vittori,^{22a,22b} I. Vivarelli,¹⁵¹ S. Vlachos,¹⁰ M. Vogel,¹⁷⁸ P. Vokac,¹³⁰ G. Volpi,^{126a,126b} H. von der Schmitt,¹⁰³ E. von Toerne,²³ V. Vorobel,¹³¹ K. Vorobev,¹⁰⁰ M. Vos,¹⁷⁰ R. Voss,³² J. H. Vossebeld,⁷⁷ N. Vranjes,¹⁴ M. Vranjes Milosavljevic,¹⁴ V. Vrba,¹³⁰ M. Vreeswijk,¹⁰⁹ R. Vuillermet,³² I. Vukotic,³³

P. Wagner,²³ W. Wagner,¹⁷⁸ J. Wagner-Kuhr,¹⁰² H. Wahlberg,⁷⁴ S. Wahrenmund,⁴⁷ J. Wakabayashi,¹⁰⁵ J. Walder,⁷⁵ R. Walker,¹⁰² W. Walkowiak,¹⁴³ V. Wallangen,^{148a,148b} C. Wang,^{35b} C. Wang,^{36b,ww} F. Wang,¹⁷⁶ H. Wang,¹⁶ H. Wang,³ J. Wang,⁴⁵ J. Wang,¹⁵² Q. Wang,¹¹⁵ R. Wang,⁶ S. M. Wang,¹⁵³ T. Wang,³⁸ W. Wang,^{153,xx} W. Wang,^{36a} Z. Wang,^{36c} C. Wanotayaroj,¹¹⁸ A. Warburton,⁹⁰ C. P. Ward,³⁰ D. R. Wardrope,⁸¹ A. Washbrook,⁴⁹ P. M. Watkins,¹⁹ A. T. Watson,¹⁹ M. F. Watson,¹⁹ G. Watts,¹⁴⁰ S. Watts,⁸⁷ B. M. Waugh,⁸¹ A. F. Webb,¹¹ S. Webb,⁸⁶ M. S. Weber,¹⁸ S. W. Weber,¹⁷⁷ S. A. Weber,³¹ J. S. Webster,⁶ A. R. Weidberg,¹²² B. Weinert,⁶⁴ J. Weingarten,⁵⁷ M. Weirich,⁸⁶ C. Weiser,⁵¹ H. Weits,¹⁰⁹ P. S. Wells,³² T. Wenaus,²⁷ T. Wengler,³² S. Wenig,³² N. Vermes,²³ M. D. Werner,⁶⁷ P. Werner,³² M. Wessels,^{60a} T. D. Weston,¹⁸ K. Whalen,¹¹⁸ N. L. Whallon,¹⁴⁰ A. M. Wharton,⁷⁵ A. S. White,⁹² A. White,⁸ M. J. White,¹ R. White,^{34b} D. Whiteson,¹⁶⁶ B. W. Whitmore,⁷⁵ F. J. Wickens,¹³³ W. Wiedenmann,¹⁷⁶ M. Wielers,¹³³ C. Wiglesworth,³⁹ L. A. M. Wiik-Fuchs,⁵¹ A. Wildauer,¹⁰³ F. Wilk,⁸⁷ H. G. Wilkens,³² H. H. Williams,¹²⁴ S. Williams,¹⁰⁹ C. Willis,⁹³ S. Willocq,⁸⁹ J. A. Wilson,¹⁹ I. Wingerter-Seez,⁵ E. Winkels,¹⁵¹ F. Winklmeier,¹¹⁸ O. J. Winston,¹⁵¹ B. T. Winter,²³ M. Wittgen,¹⁴⁵ M. Wobisch,^{82,v} T. M. H. Wolf,¹⁰⁹ R. Wolff,⁸⁸ M. W. Wolter,⁴² H. Wolters,^{128a,128c} V. W. S. Wong,¹⁷¹ S. D. Worm,¹⁹ B. K. Wosiek,⁴² J. Wotschack,³² K. W. Wozniak,⁴² M. Wu,³³ S. L. Wu,¹⁷⁶ X. Wu,⁵² Y. Wu,⁹² T. R. Wyatt,⁸⁷ B. M. Wynne,⁴⁹ S. Xella,³⁹ Z. Xi,⁹² L. Xia,^{35c} D. Xu,^{35a} L. Xu,²⁷ T. Xu,¹³⁸ B. Yabsley,¹⁵² S. Yacoob,^{147a} D. Yamaguchi,¹⁵⁹ Y. Yamaguchi,¹²⁰ A. Yamamoto,⁶⁹ S. Yamamoto,¹⁵⁷ T. Yamanaka,¹⁵⁷ M. Yamatani,¹⁵⁷ K. Yamauchi,¹⁰⁵ Y. Yamazaki,⁷⁰ Z. Yan,²⁴ H. Yang,^{36c} H. Yang,¹⁶ Y. Yang,¹⁵³ Z. Yang,¹⁵ W-M. Yao,¹⁶ Y. C. Yap,⁸³ Y. Yasu,⁶⁹ E. Yatsenko,⁵ K. H. Yau Wong,²³ J. Ye,⁴³ S. Ye,²⁷ I. Yeletsikh,⁶⁸ E. Yigitbasi,²⁴ E. Yildirim,⁸⁶ K. Yorita,¹⁷⁴ K. Yoshihara,¹²⁴ C. Young,¹⁴⁵ C. J. S. Young,³² J. Yu,⁸ J. Yu,⁶⁷ S. P. Y. Yuen,²³ I. Yusuff,^{30,y} B. Zabinski,⁴² G. Zacharis,¹⁰ R. Zaidan,¹³ A. M. Zaitsev,^{132,mm} N. Zakharchuk,⁴⁵ J. Zalieckas,¹⁵ A. Zaman,¹⁵⁰ S. Zambito,⁵⁹ D. Zanzi,⁹¹ C. Zeitnitz,¹⁷⁸ G. Zemaityte,¹²² A. Zemla,^{41a} J. C. Zeng,¹⁶⁹ Q. Zeng,¹⁴⁵ O. Zenin,¹³² T. Ženiš,^{146a} D. Zerwas,¹¹⁹ D. Zhang,⁹² F. Zhang,¹⁷⁶ G. Zhang,^{36a,zz} H. Zhang,^{35b} J. Zhang,⁶ L. Zhang,⁵¹ L. Zhang,^{36a} M. Zhang,¹⁶⁹ P. Zhang,^{35b} R. Zhang,²³ R. Zhang,^{36a,ww} X. Zhang,^{36b} Y. Zhang,^{35a} Z. Zhang,¹¹⁹ X. Zhao,⁴³ Y. Zhao,^{36b,aaa} Z. Zhao,^{36a} A. Zhemchugov,⁶⁸ B. Zhou,⁹² C. Zhou,¹⁷⁶ L. Zhou,⁴³ M. Zhou,^{35a} M. Zhou,¹⁵⁰ N. Zhou,^{35c} C. G. Zhu,^{36b} H. Zhu,^{35a} J. Zhu,⁹² Y. Zhu,^{36a} X. Zhuang,^{35a} K. Zhukov,⁹⁸ A. Zibell,¹⁷⁷ D. Zieminska,⁶⁴ N. I. Zimine,⁶⁸ C. Zimmermann,⁸⁶ S. Zimmermann,⁵¹ Z. Zinonos,¹⁰³ M. Zinser,⁸⁶ M. Ziolkowski,¹⁴³ L. Živković,¹⁴ G. Zobernig,¹⁷⁶ A. Zoccoli,^{22a,22b} R. Zou,³³ M. zur Nedden,¹⁷ and L. Zwalinski³²

(ATLAS Collaboration)

¹*Department of Physics, University of Adelaide, Adelaide, Australia*²*Physics Department, SUNY Albany, Albany New York, USA*³*Department of Physics, University of Alberta, Edmonton Alberta, Canada*^{4a}*Department of Physics, Ankara University, Ankara, Turkey*^{4b}*Istanbul Aydin University, Istanbul, Turkey*^{4c}*Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey*⁵*LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France*⁶*High Energy Physics Division, Argonne National Laboratory, Argonne Illinois, USA*⁷*Department of Physics, University of Arizona, Tucson Arizona, USA*⁸*Department of Physics, The University of Texas at Arlington, Arlington Texas, USA*⁹*Physics Department, National and Kapodistrian University of Athens, Athens, Greece*¹⁰*Physics Department, National Technical University of Athens, Zografou, Greece*¹¹*Department of Physics, The University of Texas at Austin, Austin Texas, USA*¹²*Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*¹³*Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain*¹⁴*Institute of Physics, University of Belgrade, Belgrade, Serbia*¹⁵*Department for Physics and Technology, University of Bergen, Bergen, Norway*¹⁶*Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley California, USA*¹⁷*Department of Physics, Humboldt University, Berlin, Germany*¹⁸*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*¹⁹*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*^{20a}*Department of Physics, Bogazici University, Istanbul, Turkey*^{20b}*Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey*^{20d}*Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey*^{20c}*Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey*

- ²¹*Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
- ^{22a}*INFN Sezione di Bologna, Italy*
- ^{22b}*Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy*
- ²³*Physikalisches Institut, University of Bonn, Bonn, Germany*
- ²⁴*Department of Physics, Boston University, Boston Massachusetts, USA*
- ²⁵*Department of Physics, Brandeis University, Waltham Massachusetts, USA*
- ^{26a}*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- ^{26b}*Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- ^{26c}*Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil*
- ^{26d}*Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil*
- ²⁷*Physics Department, Brookhaven National Laboratory, Upton New York, USA*
- ^{28a}*Transilvania University of Brasov, Brasov, Romania*
- ^{28b}*Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania*
- ^{28c}*Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania*
- ^{28d}*National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania*
- ^{28e}*University Politehnica Bucharest, Bucharest, Romania*
- ^{28f}*West University in Timisoara, Timisoara, Romania*
- ²⁹*Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- ³⁰*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ³¹*Department of Physics, Carleton University, Ottawa Ontario, Canada*
- ³²*CERN, Geneva, Switzerland*
- ³³*Enrico Fermi Institute, University of Chicago, Chicago Illinois, USA*
- ^{34a}*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- ^{34b}*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ^{35a}*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*
- ^{35b}*Department of Physics, Nanjing University, Jiangsu, China*
- ^{35c}*Physics Department, Tsinghua University, Beijing 100084, China*
- ^{36a}*Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Anhui, China*
- ^{36b}*School of Physics, Shandong University, Shandong, China*
- ^{36c}*Department of Physics and Astronomy, Key Laboratory for Particle Physics, Astrophysics and Cosmology, Ministry of Education; Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai(also at PKU-CHEP), China*
- ³⁷*Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France*
- ³⁸*Nevis Laboratory, Columbia University, Irvington New York, USA*
- ³⁹*Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark*
- ^{40a}*INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy*
- ^{40b}*Dipartimento di Fisica, Università della Calabria, Rende, Italy*
- ^{41a}*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- ^{41b}*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- ⁴²*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- ⁴³*Physics Department, Southern Methodist University, Dallas Texas, USA*
- ⁴⁴*Physics Department, University of Texas at Dallas, Richardson Texas, USA*
- ⁴⁵*DESY, Hamburg and Zeuthen, Germany*
- ⁴⁶*Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*
- ⁴⁷*Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany*
- ⁴⁸*Department of Physics, Duke University, Durham North Carolina, USA*
- ⁴⁹*SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁵⁰*INFN e Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁵¹*Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany*
- ⁵²*Departement de Physique Nucleaire et Corpusculaire, Université de Genève, Geneva, Switzerland*
- ^{53a}*INFN Sezione di Genova, Italy*
- ^{53b}*Dipartimento di Fisica, Università di Genova, Genova, Italy*
- ^{54a}*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
- ^{54b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
- ⁵⁵*II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
- ⁵⁶*SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- ⁵⁷*II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany*
- ⁵⁸*Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France*
- ⁵⁹*Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge Massachusetts, USA*

- ^{60a}*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
^{60b}*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
⁶¹*Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*
^{62a}*Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China*
^{62b}*Department of Physics, The University of Hong Kong, Hong Kong, China*
^{62c}*Department of Physics and Institute for Advanced Study, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
⁶³*Department of Physics, National Tsing Hua University, Taiwan, Taiwan*
⁶⁴*Department of Physics, Indiana University, Bloomington Indiana, USA*
⁶⁵*Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria*
⁶⁶*University of Iowa, Iowa City Iowa, USA*
⁶⁷*Department of Physics and Astronomy, Iowa State University, Ames Iowa, USA*
⁶⁸*Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia*
⁶⁹*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
⁷⁰*Graduate School of Science, Kobe University, Kobe, Japan*
⁷¹*Faculty of Science, Kyoto University, Kyoto, Japan*
⁷²*Kyoto University of Education, Kyoto, Japan*
⁷³*Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan*
⁷⁴*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
⁷⁵*Physics Department, Lancaster University, Lancaster, United Kingdom*
^{76a}*INFN Sezione di Lecce, Italy*
^{76b}*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
⁷⁷*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
⁷⁸*Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia*
⁷⁹*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
⁸⁰*Department of Physics, Royal Holloway University of London, Surrey, United Kingdom*
⁸¹*Department of Physics and Astronomy, University College London, London, United Kingdom*
⁸²*Louisiana Tech University, Ruston Louisiana, USA*
⁸³*Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France*
⁸⁴*Fysiska institutionen, Lunds universitet, Lund, Sweden*
⁸⁵*Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain*
⁸⁶*Institut für Physik, Universität Mainz, Mainz, Germany*
⁸⁷*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
⁸⁸*CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France*
⁸⁹*Department of Physics, University of Massachusetts, Amherst Massachusetts, USA*
⁹⁰*Department of Physics, McGill University, Montreal Québec, Canada*
⁹¹*School of Physics, University of Melbourne, Victoria, Australia*
⁹²*Department of Physics, The University of Michigan, Ann Arbor Michigan, USA*
⁹³*Department of Physics and Astronomy, Michigan State University, East Lansing Michigan, USA*
^{94a}*INFN Sezione di Milano, Italy*
^{94b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
⁹⁵*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus*
⁹⁶*Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Republic of Belarus*
⁹⁷*Group of Particle Physics, University of Montreal, Montreal Québec Canada*
⁹⁸*P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia*
⁹⁹*Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia*
¹⁰⁰*National Research Nuclear University MEPhI, Moscow, Russia*
¹⁰¹*D.V. Skobel'syn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia*
¹⁰²*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
¹⁰³*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
¹⁰⁴*Nagasaki Institute of Applied Science, Nagasaki, Japan*
¹⁰⁵*Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan*
^{106a}*INFN Sezione di Napoli, Italy*
^{106b}*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
¹⁰⁷*Department of Physics and Astronomy, University of New Mexico, Albuquerque New Mexico, USA*
¹⁰⁸*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands*
¹⁰⁹*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
¹¹⁰*Department of Physics, Northern Illinois University, DeKalb Illinois, USA*
¹¹¹*Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia*

- ¹¹²*Department of Physics, New York University, New York New York, USA*
- ¹¹³*Ohio State University, Columbus Ohio, USA*
- ¹¹⁴*Faculty of Science, Okayama University, Okayama, Japan*
- ¹¹⁵*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman Oklahoma, USA*
- ¹¹⁶*Department of Physics, Oklahoma State University, Stillwater Oklahoma, USA*
- ¹¹⁷*Palacký University, RCPTM, Olomouc, Czech Republic*
- ¹¹⁸*Center for High Energy Physics, University of Oregon, Eugene Oregon, USA*
- ¹¹⁹*LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France*
- ¹²⁰*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹²¹*Department of Physics, University of Oslo, Oslo, Norway*
- ¹²²*Department of Physics, Oxford University, Oxford, United Kingdom*
- ^{123a}*INFN Sezione di Pavia, Italy*
- ^{123b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- ¹²⁴*Department of Physics, University of Pennsylvania, Philadelphia Pennsylvania, USA*
- ¹²⁵*National Research Centre “Kurchatov Institute” B.P.Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia*
- ^{126a}*INFN Sezione di Pisa, Italy*
- ^{126b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ¹²⁷*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh Pennsylvania, USA*
- ^{128a}*Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa, Portugal*
- ^{128b}*Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal*
- ^{128c}*Department of Physics, University of Coimbra, Coimbra, Portugal*
- ^{128d}*Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal*
- ^{128e}*Departamento de Física, Universidade do Minho, Braga, Portugal*
- ^{128f}*Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada, Portugal*
- ^{128g}*Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal*
- ¹²⁹*Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic*
- ¹³⁰*Czech Technical University in Prague, Praha, Czech Republic*
- ¹³¹*Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic*
- ¹³²*State Research Center Institute for High Energy Physics (Protvino), NRC KI, Russia*
- ¹³³*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ^{134a}*INFN Sezione di Roma, Italy*
- ^{134b}*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
- ^{135a}*INFN Sezione di Roma Tor Vergata, Italy*
- ^{135b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{136a}*INFN Sezione di Roma Tre, Italy*
- ^{136b}*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
- ^{137a}*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca, Morocco*
- ^{137b}*Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat, Morocco*
- ^{137c}*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- ^{137d}*Faculté des Sciences, Université Mohamed Premier and LTPM, Oujda, Morocco*
- ^{137e}*Faculté des sciences, Université Mohammed V, Rabat, Morocco*
- ¹³⁸*DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France*
- ¹³⁹*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz California, USA*
- ¹⁴⁰*Department of Physics, University of Washington, Seattle Washington, USA*
- ¹⁴¹*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹⁴²*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁴³*Department Physik, Universität Siegen, Siegen, Germany*
- ¹⁴⁴*Department of Physics, Simon Fraser University, Burnaby British Columbia, Canada*
- ¹⁴⁵*SLAC National Accelerator Laboratory, Stanford California, USA*
- ^{146a}*Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic*
- ^{146b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*
- ^{147a}*Department of Physics, University of Cape Town, Cape Town, South Africa*
- ^{147b}*Department of Physics, University of Johannesburg, Johannesburg, South Africa*
- ^{147c}*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
- ^{148a}*Department of Physics, Stockholm University, Sweden*
- ^{148b}*The Oskar Klein Centre, Stockholm, Sweden*
- ¹⁴⁹*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
- ¹⁵⁰*Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook New York, USA*

- ¹⁵¹*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
- ¹⁵²*School of Physics, University of Sydney, Sydney, Australia*
- ¹⁵³*Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ¹⁵⁴*Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel*
- ¹⁵⁵*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
- ¹⁵⁶*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
- ¹⁵⁷*International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
- ¹⁵⁸*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
- ¹⁵⁹*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
- ¹⁶⁰*Tomsk State University, Tomsk, Russia*
- ¹⁶¹*Department of Physics, University of Toronto, Toronto Ontario, Canada*
- ^{162a}*INFN-TIFPA, Italy*
- ^{162b}*University of Trento, Trento, Italy*
- ^{163a}*TRIUMF, Vancouver British Columbia, Canada*
- ^{163b}*Department of Physics and Astronomy, York University, Toronto Ontario, Canada*
- ¹⁶⁴*Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan*
- ¹⁶⁵*Department of Physics and Astronomy, Tufts University, Medford Massachusetts, USA*
- ¹⁶⁶*Department of Physics and Astronomy, University of California Irvine, Irvine California, USA*
- ^{167a}*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
- ^{167b}*ICTP, Trieste, Italy*
- ^{167c}*Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
- ¹⁶⁸*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
- ¹⁶⁹*Department of Physics, University of Illinois, Urbana Illinois, USA*
- ¹⁷⁰*Instituto de Fisica Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Spain*
- ¹⁷¹*Department of Physics, University of British Columbia, Vancouver British Columbia, Canada*
- ¹⁷²*Department of Physics and Astronomy, University of Victoria, Victoria British Columbia, Canada*
- ¹⁷³*Department of Physics, University of Warwick, Coventry, United Kingdom*
- ¹⁷⁴*Waseda University, Tokyo, Japan*
- ¹⁷⁵*Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
- ¹⁷⁶*Department of Physics, University of Wisconsin, Madison Wisconsin, USA*
- ¹⁷⁷*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
- ¹⁷⁸*Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
- ¹⁷⁹*Department of Physics, Yale University, New Haven Connecticut, USA*
- ¹⁸⁰*Yerevan Physics Institute, Yerevan, Armenia*
- ¹⁸¹*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*
- ¹⁸²*Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan*

^aDeceased.

^bAlso at Department of Physics, King's College London, London, United Kingdom.

^cAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^dAlso at Novosibirsk State University, Novosibirsk, Russia.

^eAlso at TRIUMF, Vancouver BC, Canada.

^fAlso at Department of Physics & Astronomy, University of Louisville, Louisville, KY, USA.

^gAlso at Physics Department, An-Najah National University, Nablus, Palestine.

^hAlso at Department of Physics, California State University, Fresno CA, USA.

ⁱAlso at Department of Physics, University of Fribourg, Fribourg, Switzerland.

^jAlso at II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany.

^kAlso at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain.

^lAlso at Departamento de Física e Astronomia, Faculdade de Ciências, Universidade do Porto, Portugal.

^mAlso at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

ⁿAlso at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China.

^oAlso at Università di Napoli Parthenope, Napoli, Italy.

^pAlso at Institute of Particle Physics (IPP), Canada.

^qAlso at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania.

^rAlso at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

^sAlso at Borough of Manhattan Community College, City University of New York, New York City, USA.

^tAlso at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

^uAlso at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town, South Africa.

^vAlso at Louisiana Tech University, Ruston LA, USA.

- ^w Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.
- ^x Also at Graduate School of Science, Osaka University, Osaka, Japan.
- ^y Also at Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany.
- ^z Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands.
- ^{aa} Also at Department of Physics, The University of Texas at Austin, Austin TX, USA.
- ^{bb} Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.
- ^{cc} Also at CERN, Geneva, Switzerland.
- ^{dd} Also at Georgian Technical University (GTU), Tbilisi, Georgia.
- ^{ee} Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.
- ^{ff} Also at Manhattan College, New York NY, USA.
- ^{gg} Also at Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile.
- ^{hh} Also at Department of Physics, The University of Michigan, Ann Arbor MI, USA.
- ⁱⁱ Also at The City College of New York, New York NY, USA.
- ^{jj} Also at School of Physics, Shandong University, Shandong, China.
- ^{kk} Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada, Portugal.
- ^{ll} Also at Department of Physics, California State University, Sacramento CA, USA.
- ^{mm} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
- ⁿⁿ Also at Departement de Physique Nucleaire et Corpusculaire, Université de Genève, Geneva, Switzerland.
- ^{oo} Also at Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Barcelona, Spain.
- ^{pp} Also at School of Physics, Sun Yat-sen University, Guangzhou, China.
- ^{qq} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria.
- ^{rr} Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia.
- ^{ss} Also at National Research Nuclear University MEPhI, Moscow, Russia.
- ^{tt} Also at Department of Physics, Stanford University, Stanford CA, USA.
- ^{uu} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
- ^{vv} Also at Giresun University, Faculty of Engineering, Turkey.
- ^{ww} Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.
- ^{xx} Also at Department of Physics, Nanjing University, Jiangsu, China.
- ^{yy} Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.
- ^{zz} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
- ^{aaa} Also at LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay, France.